



TESTING GENERAL RELATIVITY WITH GRAVITATIONAL WAVES

A "NO-HAIR" TEST FOR BINARY BLACK HOLES

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on behalf of

the LIGO Scientific Collaboration and Virgo Collaboration

DISCLAIMERS

- testing the no-hair theorem doesn't mean we don't trust the theorem
- on the contrary the theorem helps to confirm if merger remnants are truly black holes
- we will assume that astrophysical black holes carry no electronic charge
- LIGO's discoveries are possibly from Kerr black holes
- in principle we can constrain the magnitude of charge

OVERVIEW

- null tests of general relativity
- tests of propagation of propagation
- test of black hole hypothesis of LIGO's detections

WHY ARE BINARY BLACK HOLES GOOD TESTBEDS FOR GENERAL RELATIVITY?

$$G_{\mu\nu} = \frac{1}{\kappa} T_{\mu\nu}, \quad \kappa = \frac{c^4}{8\pi G} \sim 5 \times 10^{49} \text{ dynes}$$

$$L \sim \frac{32\eta^2 c^5}{5G} \left(\frac{v}{c}\right)^{10}, \quad \eta = \frac{m_1 m_2}{(m_1 + m_2)^2}, \quad \frac{2c^5}{5G} \simeq 10^{59} \text{ erg s}^{-1}$$

TESTS OF GR AND ALTERNATIVES

- null tests of GR
 - assume that GR is correct and look for small deviations from GR
 - e.g. search for tails of gravitational waves
- tests of alternative theories of gravity
 - modified phase evolution of GR waveform by introducing new terms
 - could potentially arise from an alternative theory
 - e.g. massive graviton, dipole radiation, scalar modes, ...

NULL TESTS WITH LIGO'S DETECTIONS

$$\tilde{h}(f) = \mathcal{A}(f)e^{i\varphi(f)}$$

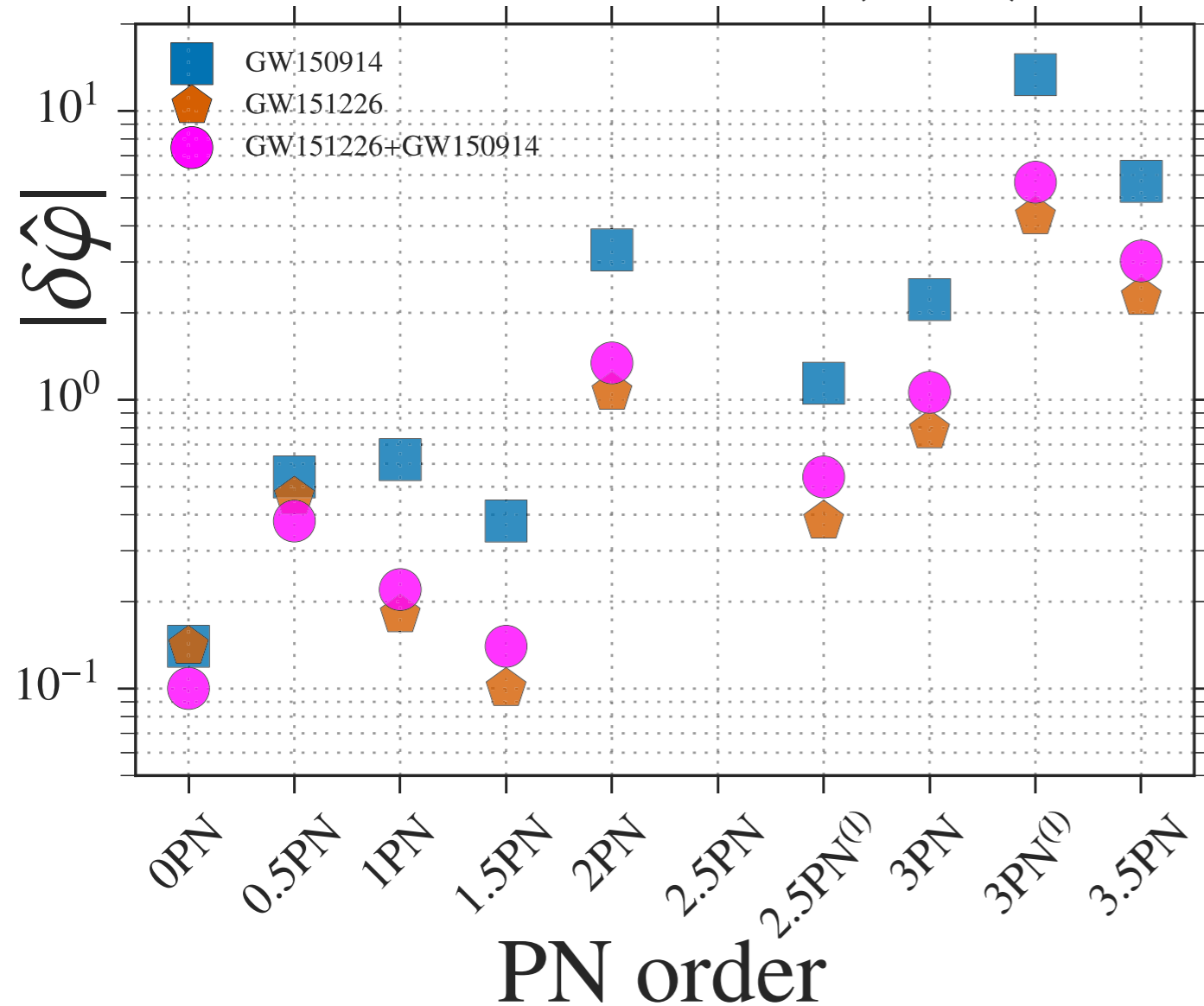
$$\begin{aligned} \varphi(f) = & \varphi_{\text{ref}} + 2\pi f t_{\text{ref}} + \varphi_{\text{Newt}}(Mf)^{-5/3} \\ & + \varphi_{0.5\text{PN}}(Mf)^{-4/3} + \varphi_{1\text{PN}}(Mf)^{-1} \\ & + \varphi_{1.5\text{PN}}(Mf)^{-2/3} + \dots \end{aligned}$$

Blanchet and BS 1995

Arun+ 2006, Mishra+ 2010,

Yunes and Pretorius 2009, Li+ 2012

Abbott+ PRX, 6, 041015 (2016)



wave tails

mass asymmetry

spin-spin coupling

spin-orbit coupling

hereditary terms

precession

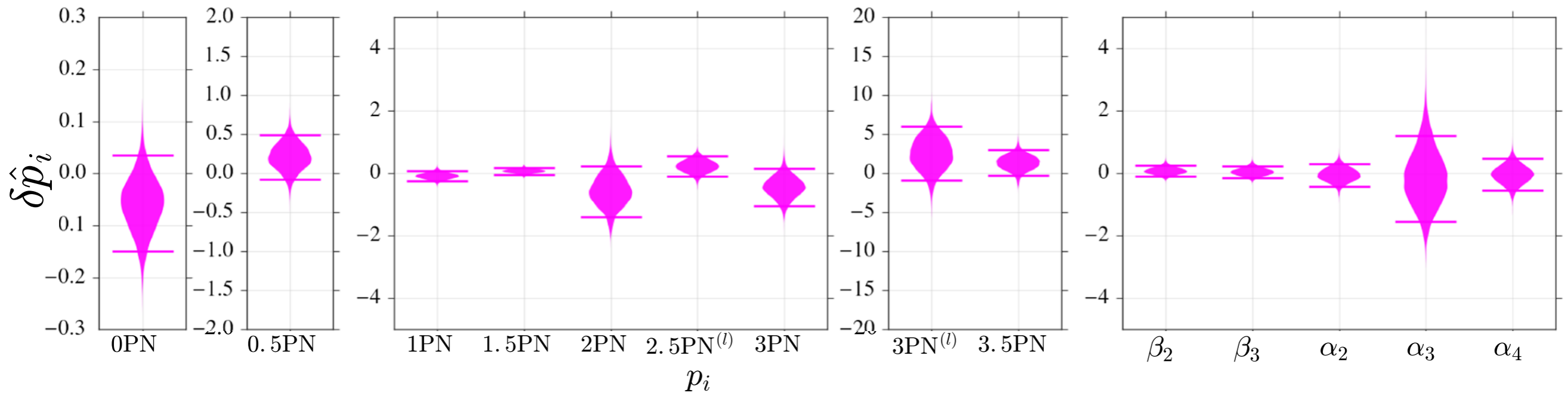
absorption of radiation by black hole

deform PN coefficients from their GR value and look for them in the data

$$\varphi_{\text{Newt}} \rightarrow \varphi_{\text{Newt}} + \delta\varphi_{\text{Newt}}$$

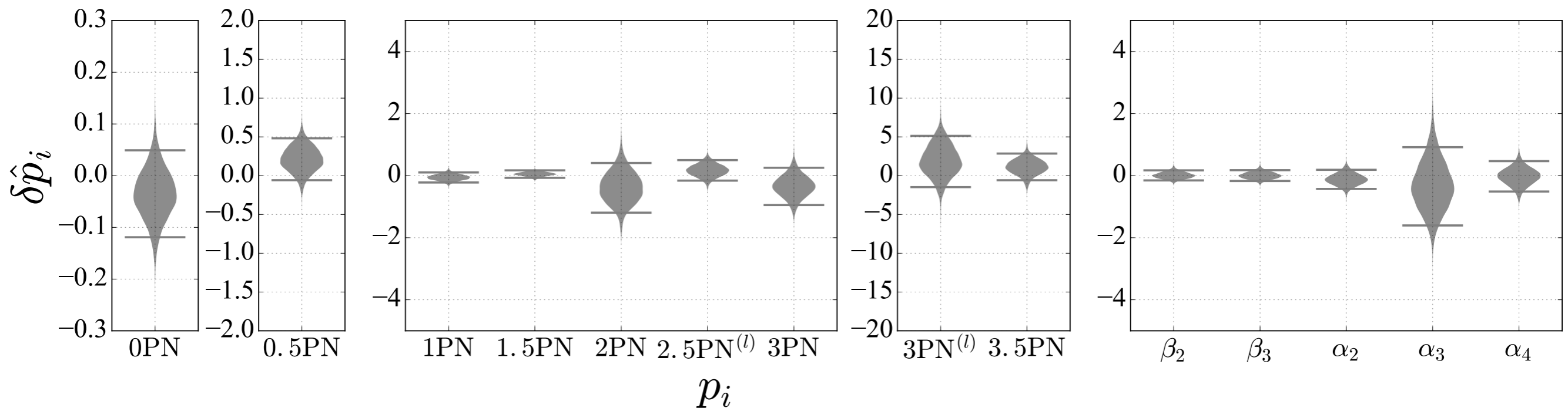
TESTS USING IMR WAVEFORM MODEL

GW150914+GW151226



Abbott+ PRX, 6, 041015 2016

GW170104



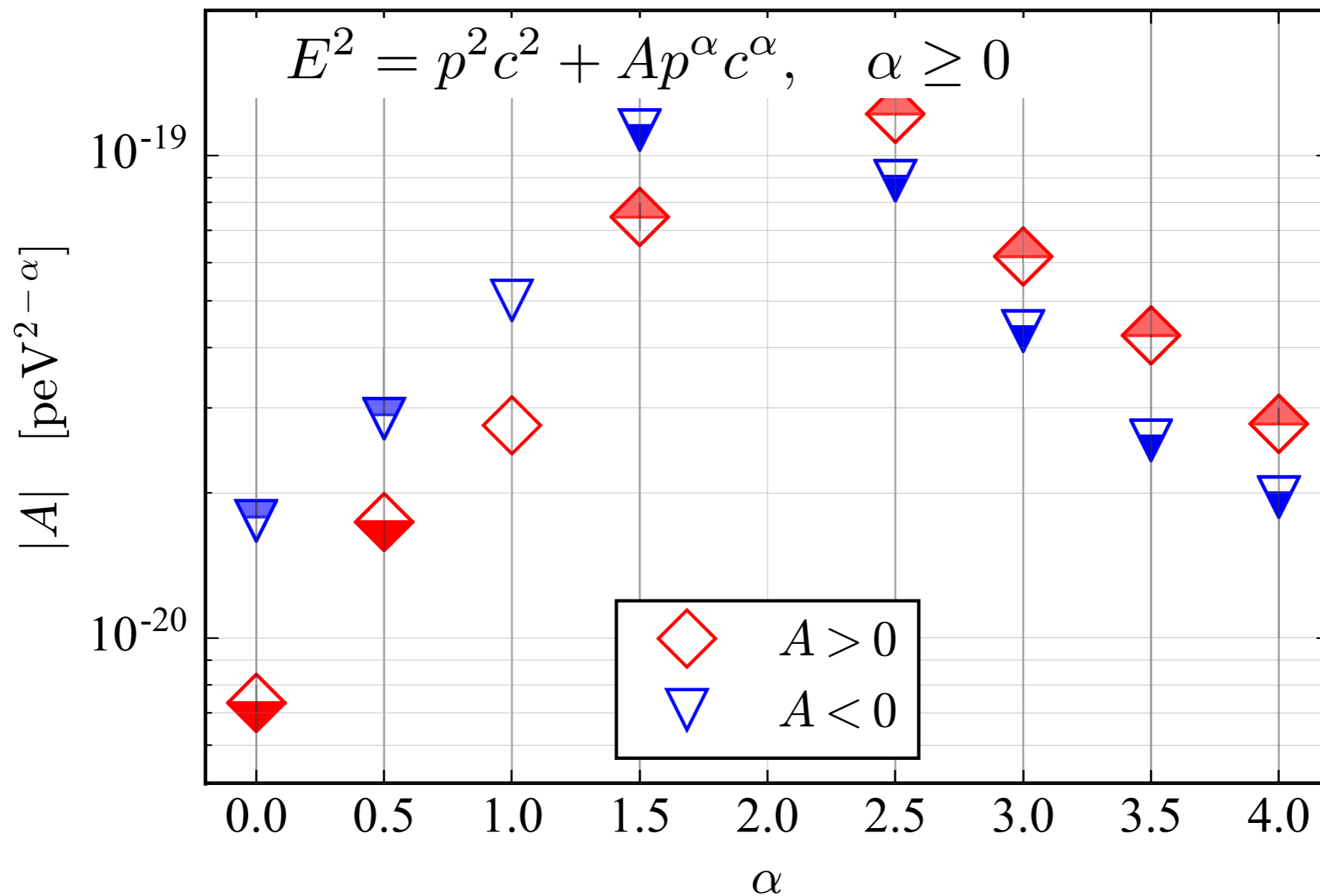
Abbott+ PRL, 118, 221101 (2017)

MODIFIED DISPERSION RELATIONS

$$E^2 = p^2 c^2 + A p^\alpha c^\alpha, \quad \alpha \geq 0$$

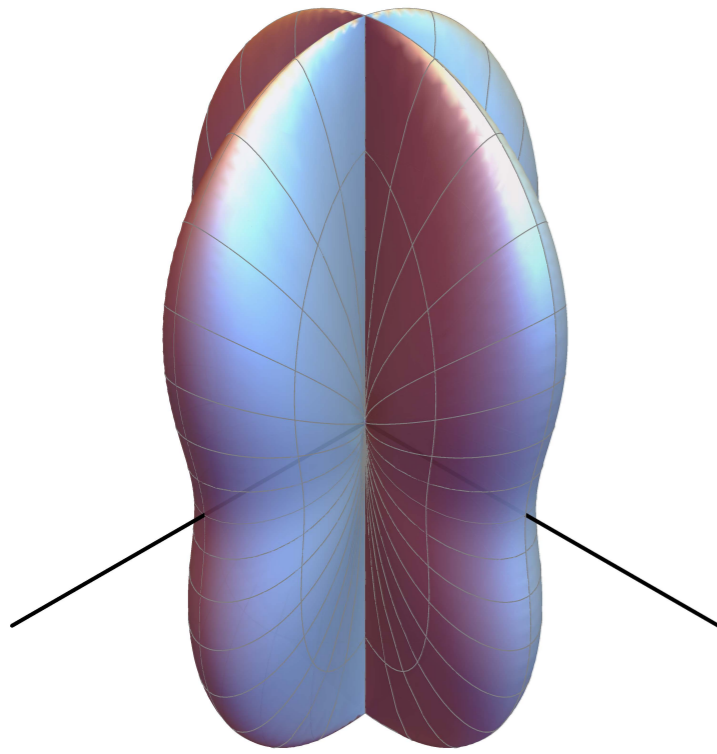
- many theories of gravity predict gravitational waves could be dispersed
- this is the first time we have tested for generic dispersion relation for gravitational waves
- first constraints in the gravity sector for superluminal gravitational waves
- find no evidence for dispersion of gravitational waves
- GW170104 bound on graviton mass: $< 7.7 \times 10^{-23}/c^2$

CONSTRAINTS ON DISPERSION OF GRAVITATIONAL WAVES

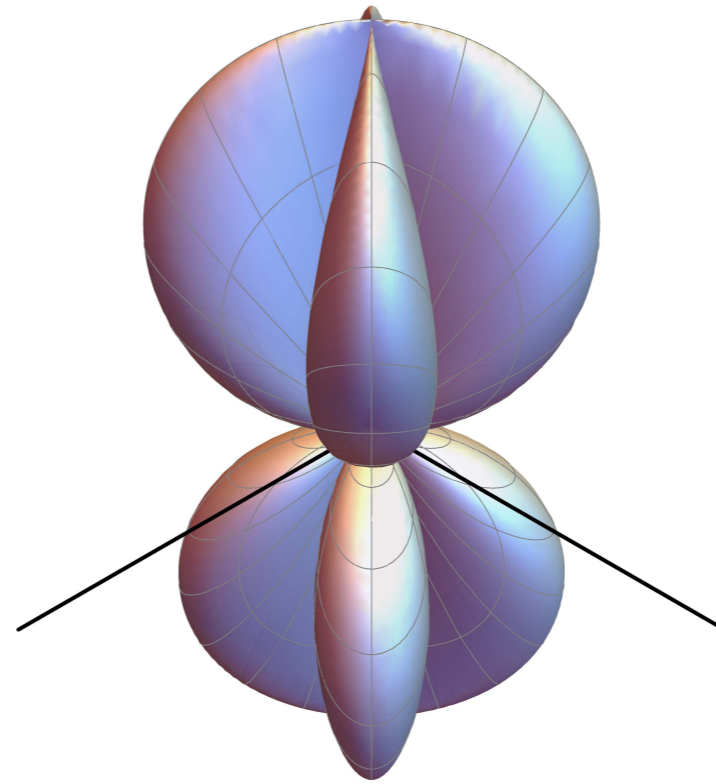


massive-graviton theories ($\alpha=0, A > 0$), multifractal spacetime ($\alpha=2.5$), doubly special relativity ($\alpha=3$), and Hor̥ava-Lifshitz and extra-dimensional theories ($\alpha=4$).

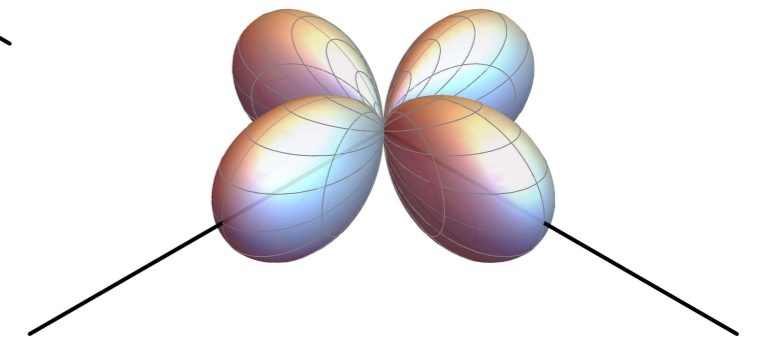
ANTENNA PATTERN FUNCTIONS DUE TO TENSOR, VECTOR AND SCALAR POLARIZATIONS



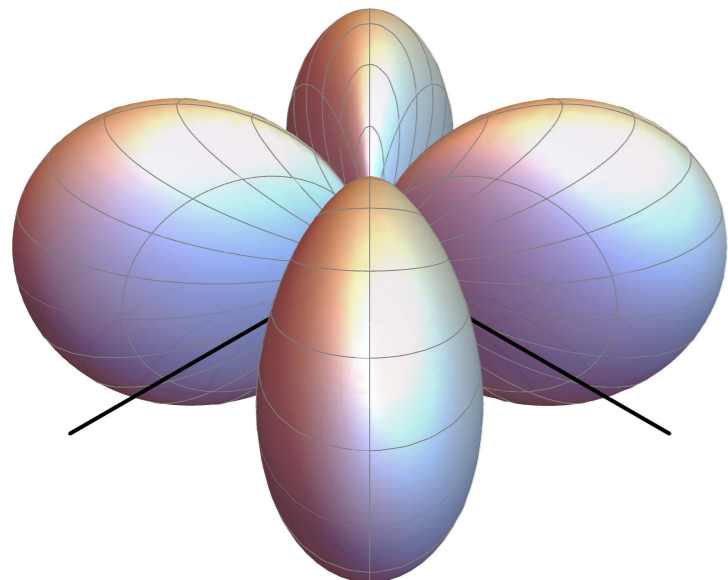
Tensor + mode



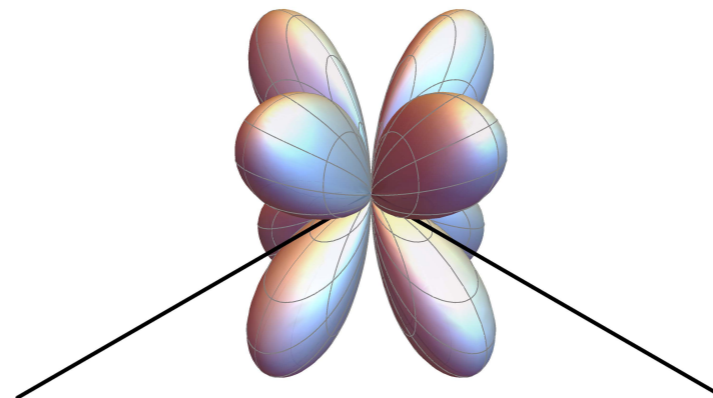
Tensor x mode



scalar mode
longitudinal and transverse



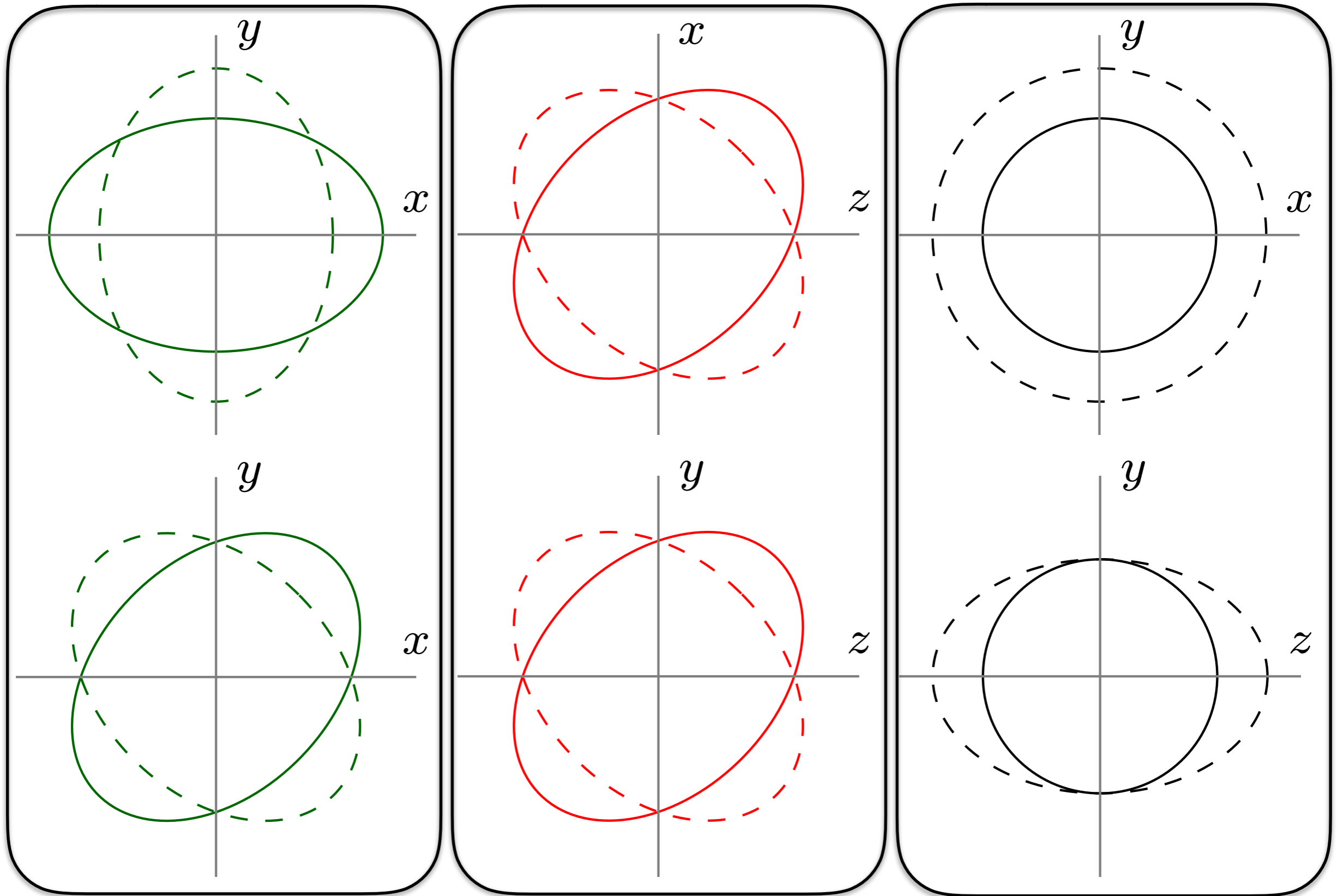
Vector x mode



Vector y mode

LVC/EPO, Max I_{si}

DEFORMATION OF A RING OF FREE PARTICLES BY DIFFERENT POLARIZATIONS



TENSOR VS SCALAR/VECTOR POLARIZATION: GW170814

- computed the Bayes factor for purely tensor modes vs purely scalar modes AND purely tensor modes vs purely vector modes
- compared to purely vector mode purely tensor mode is favored by a Bayes factor of 1:200
- compared to purely vector mode purely tensor mode is favored by a Bayes factor of 1:1,000
- 3 detectors have very poor sensitivity to a mixture model of tensor-vector-scalar modes

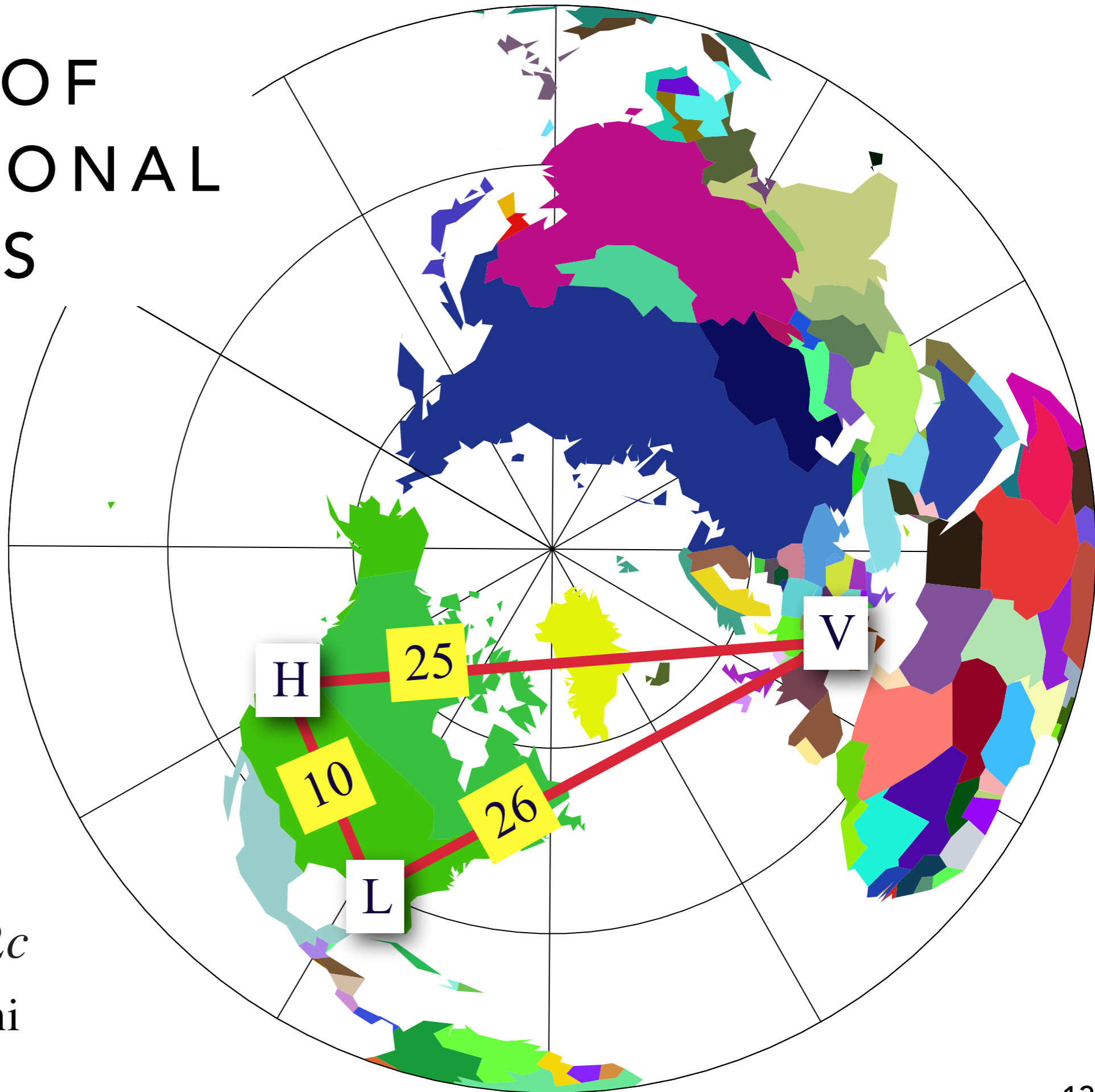
SPEED OF GRAVITATIONAL WAVES

BASELINES IN LIGHT TRAVEL TIME (MS)

USING
GW150914,
GW151226,
GW170104

$$0.55c < c_{\text{gw}} < 1.42c$$

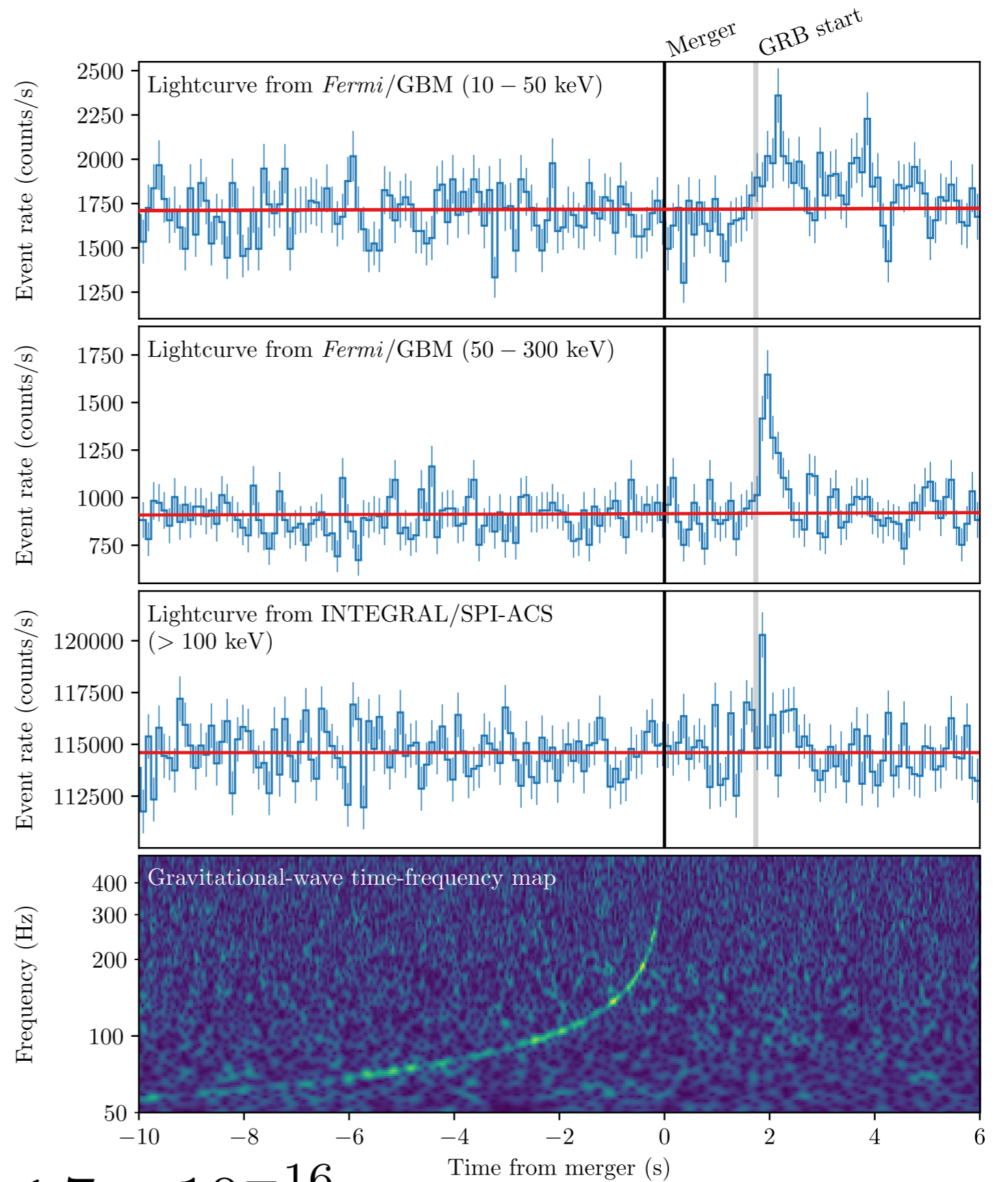
Cornish, Blas Nardini
PRL 119, 161102 (2017)



SPEED OF GRAVITATIONAL WAVES

$$\frac{1.7 \text{ s}}{40 \text{ Mpc}/c} \sim 4 \times 10^{-16}$$

$$-3 \times 10^{-15} \leq \frac{v_{\text{GW}} - v_{\text{EM}}}{v_{\text{EM}}} \leq 7 \times 10^{-16}$$



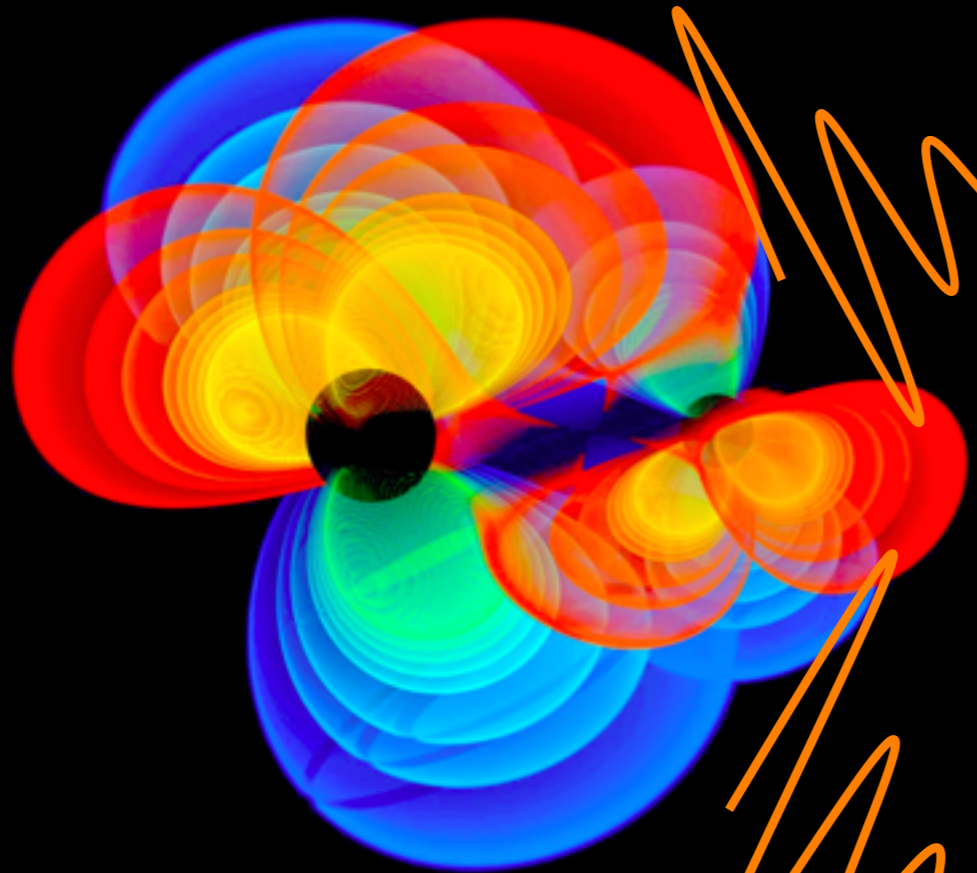
QUASI-NORMAL MODES AND NO-HAIR

- ❖ Deformed black holes emit quasi-normal modes

- ❖ complex frequencies depend only on the mass and spin

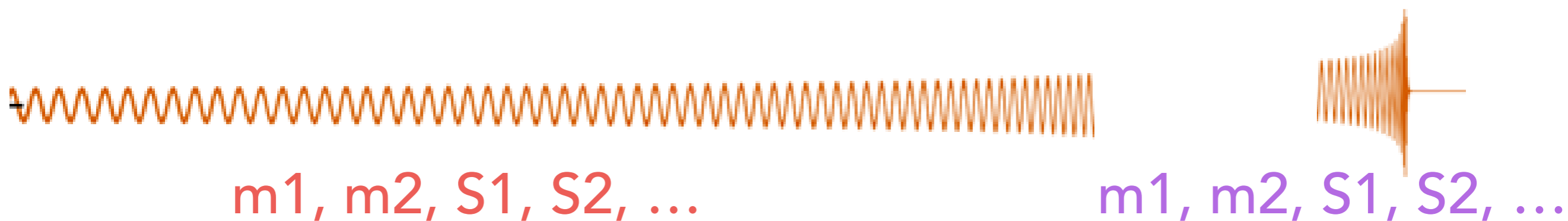
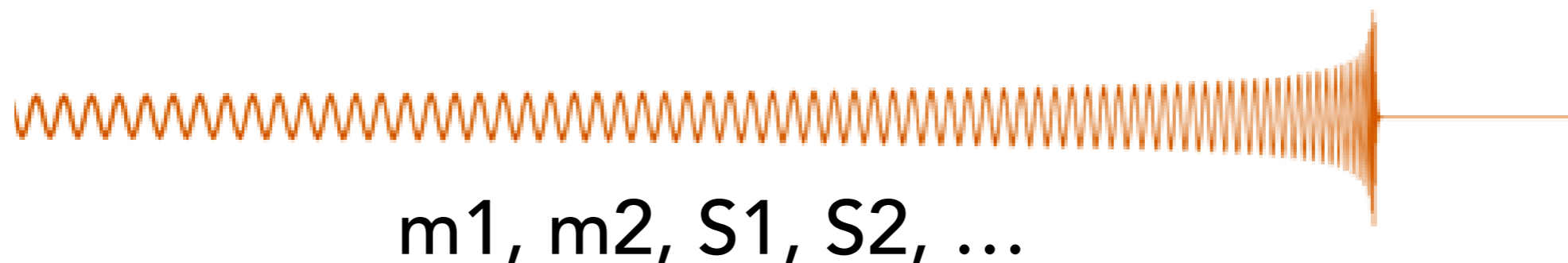
- ❖ Measuring two or modes would provide a smoking gun evidence of Kerr black holes

- ❖ If modes depend on other parameters, consistency between different mode frequencies would fail

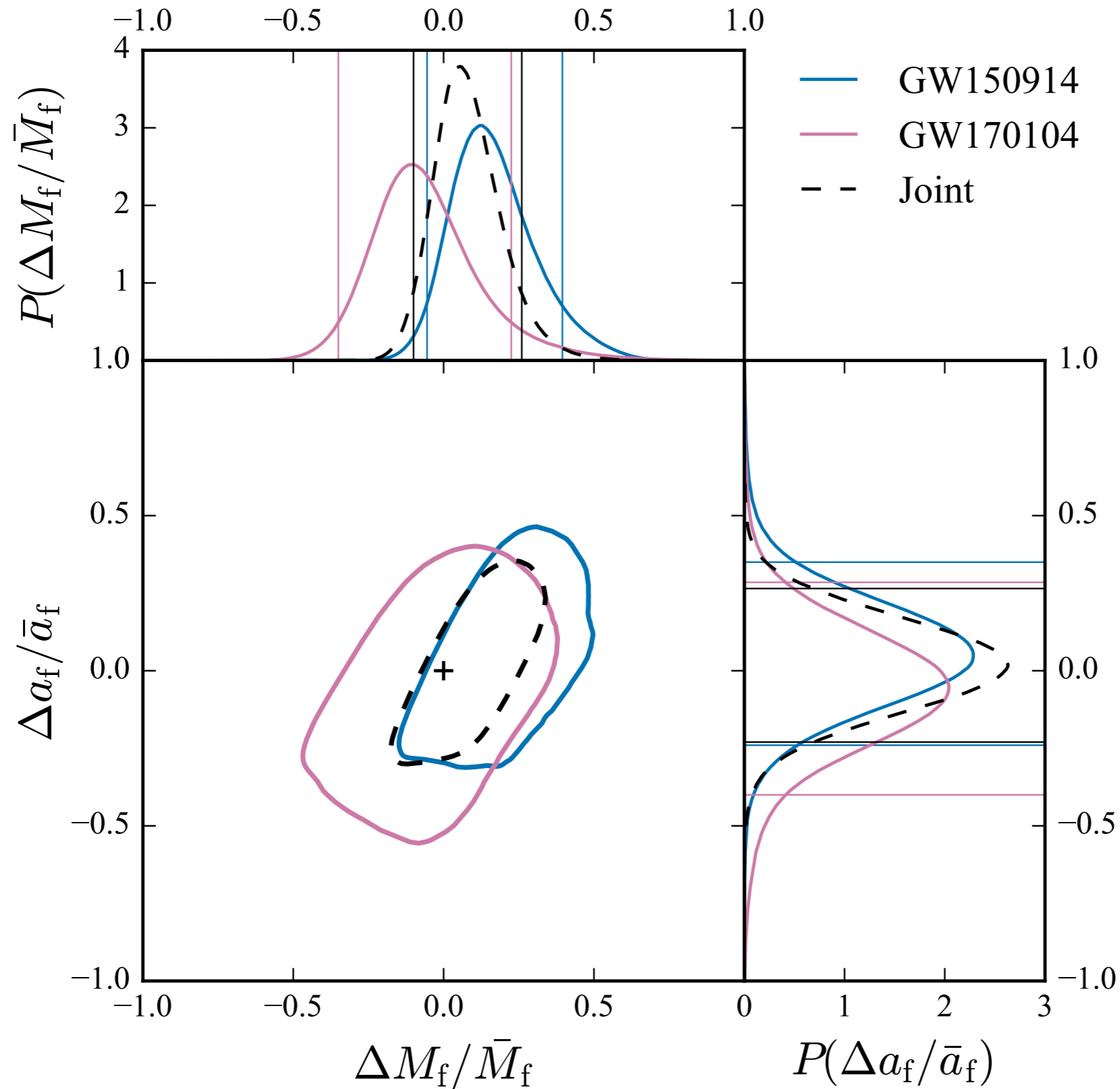


Dreyer+ 2004, Berti+ 2006, Berti+ 2007, Kamaretsos+ 2012, Gossan +2012, Bhagwat+ 2017, Brito+ 2018

IS THE SIGNAL CONSISTENT WITH A BLACK HOLE REMNANT

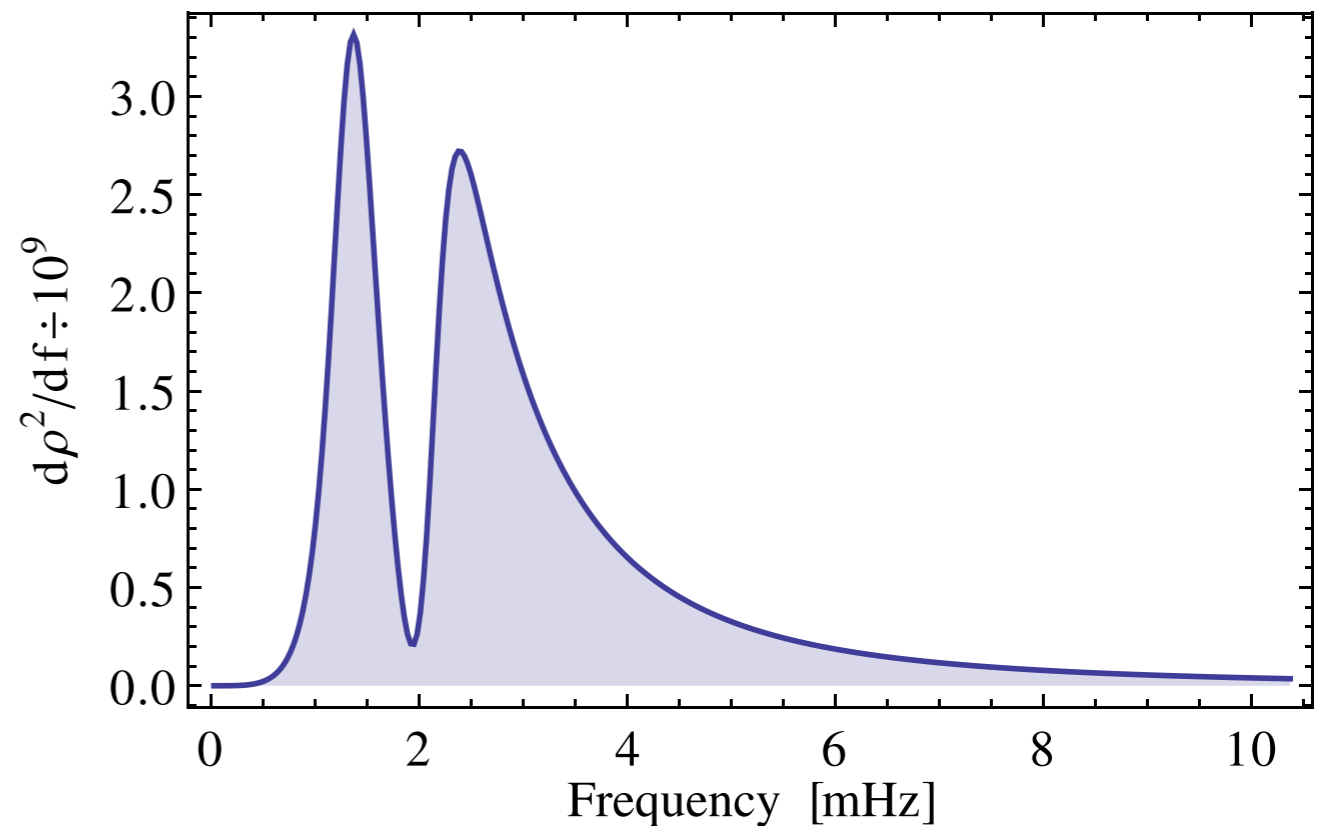
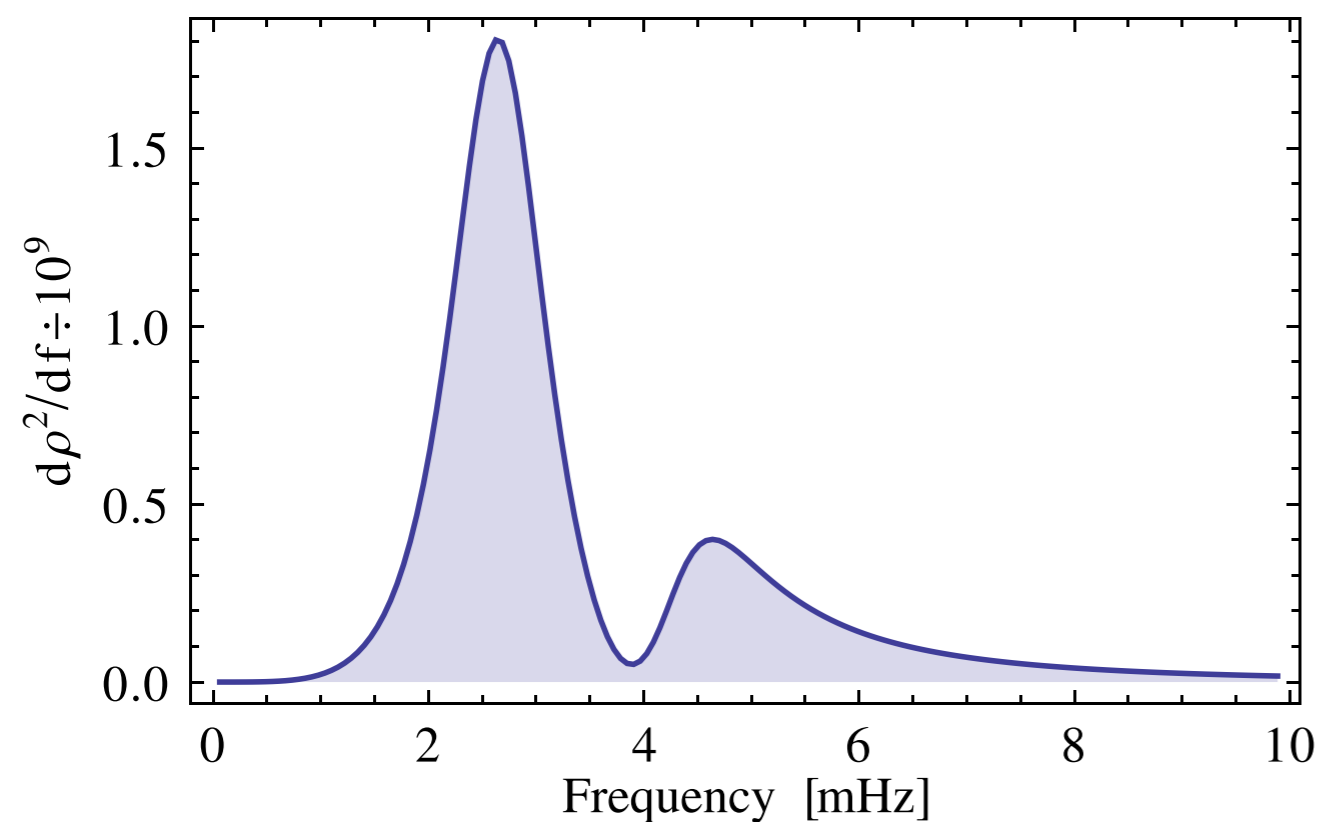


IS THE SIGNAL CONSISTENT WITH A BLACK HOLE REMNANT



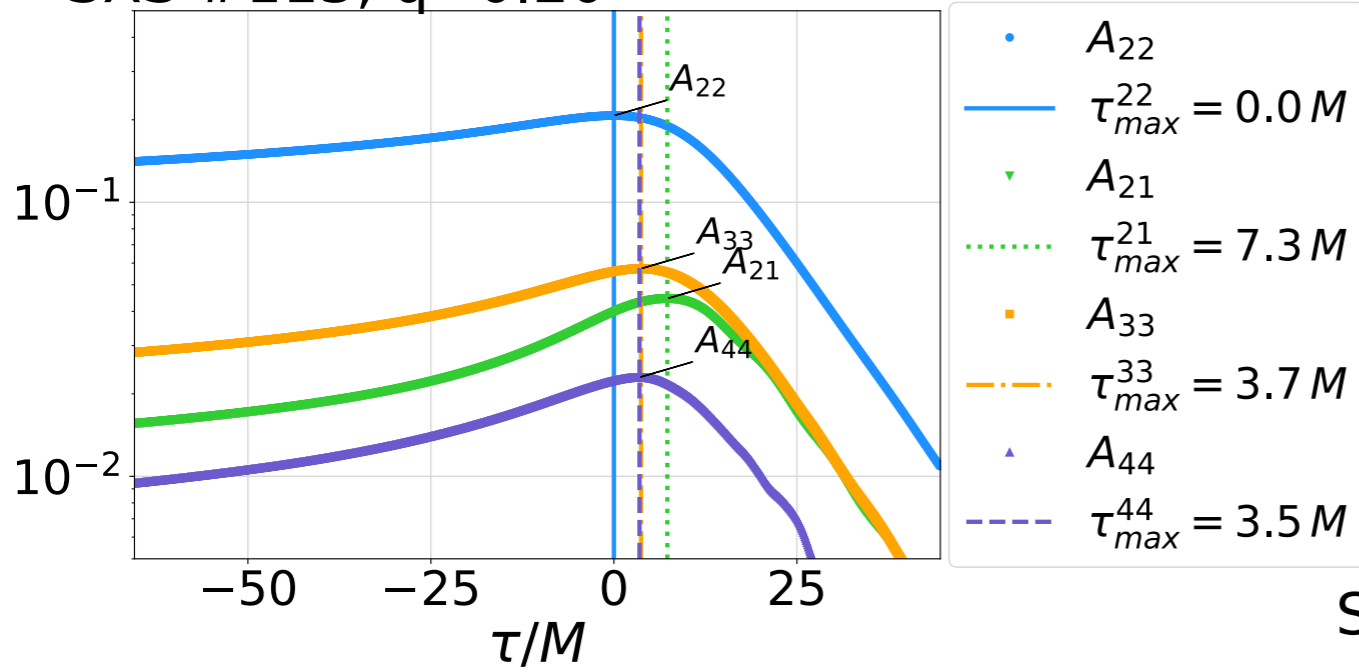
QNM SIGNATURE IN BINARY BLACK HOLES

- Quasi normal mode signal example in LISA:
 - left total mass of 5 million solar mass
 - right total mass of 50 million solar masses
- mass ratio of 1:10 in both cases

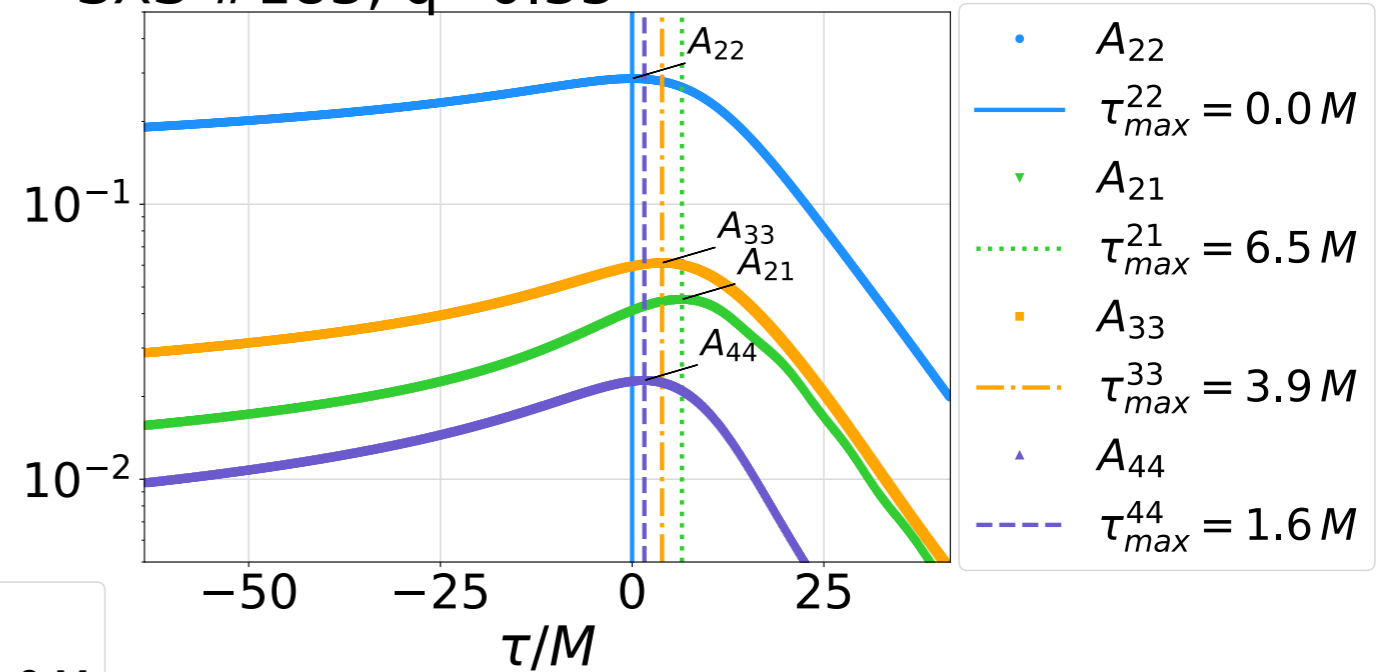


HIGHER MODES FROM INSPIRAL TO RINGDOWN

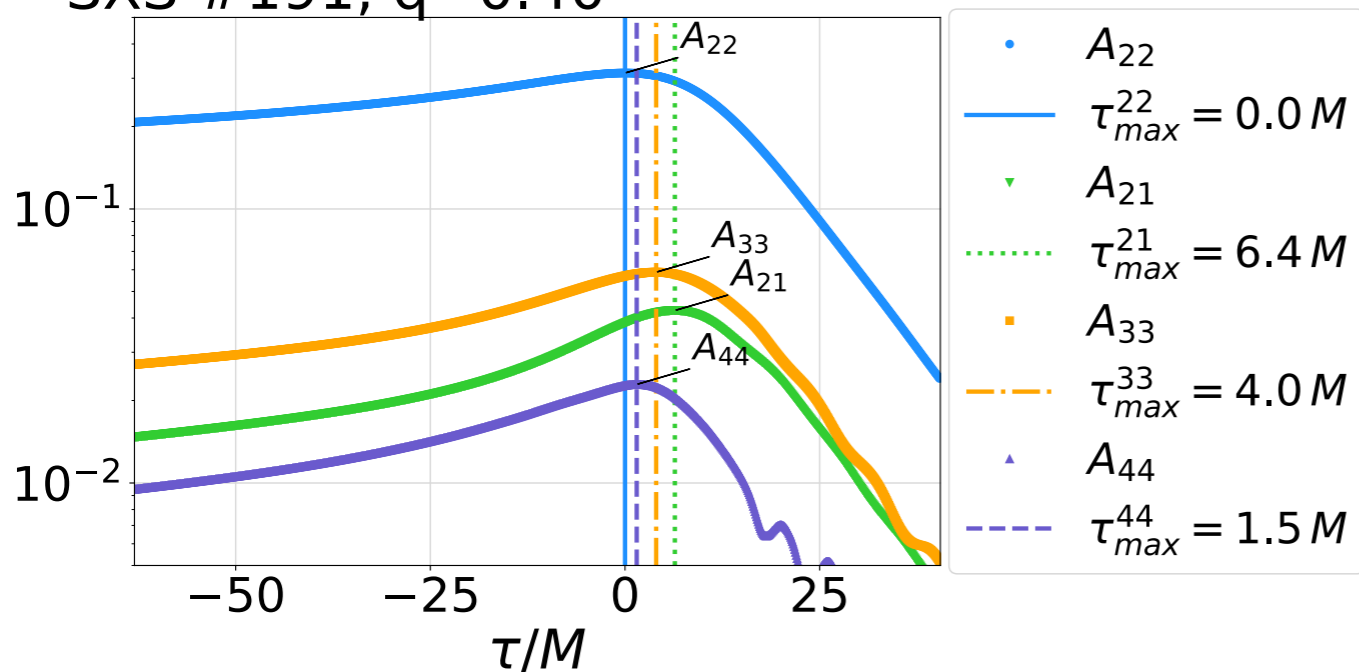
SXS #113, $q=0.20$



SXS #183, $q=0.33$

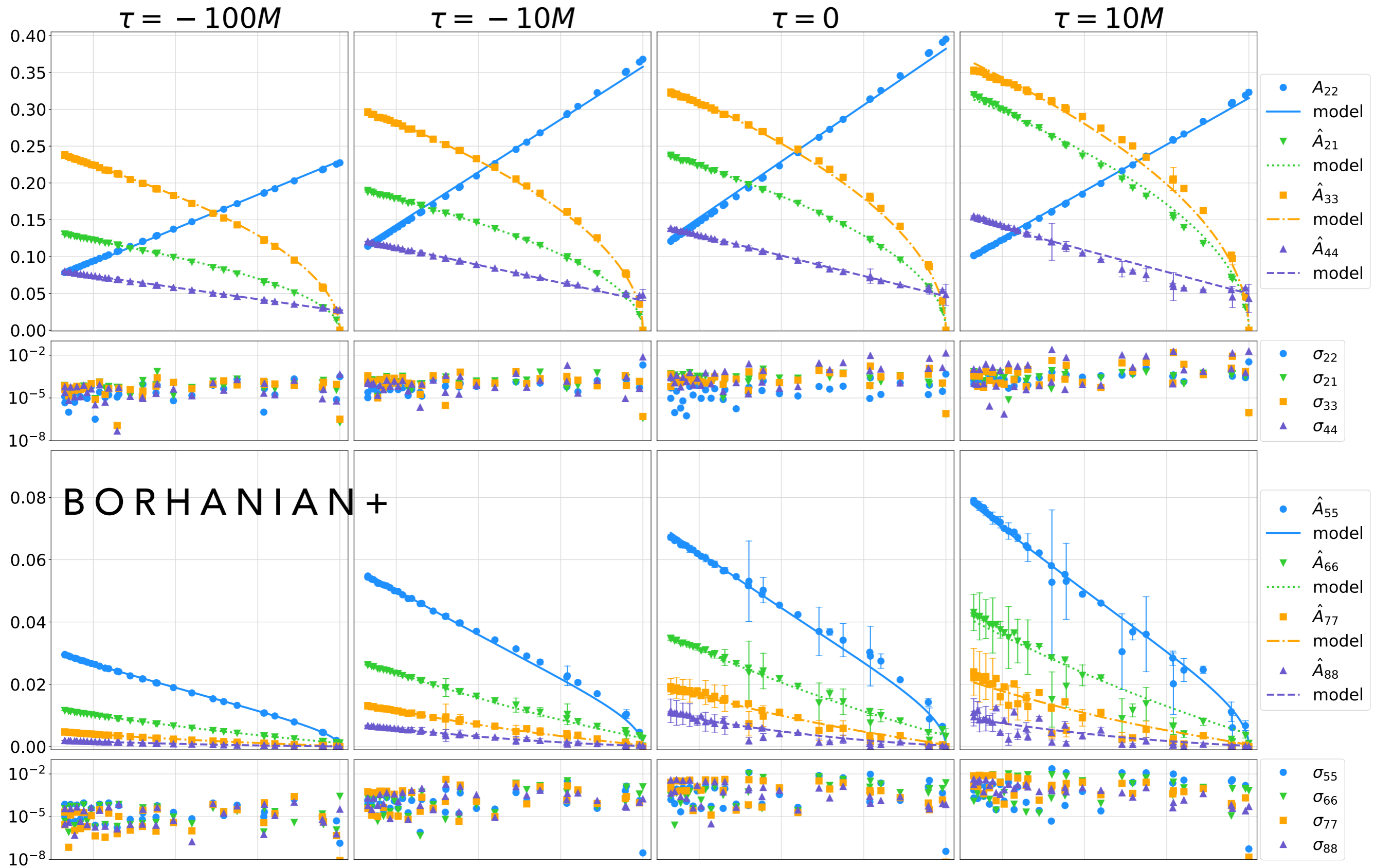


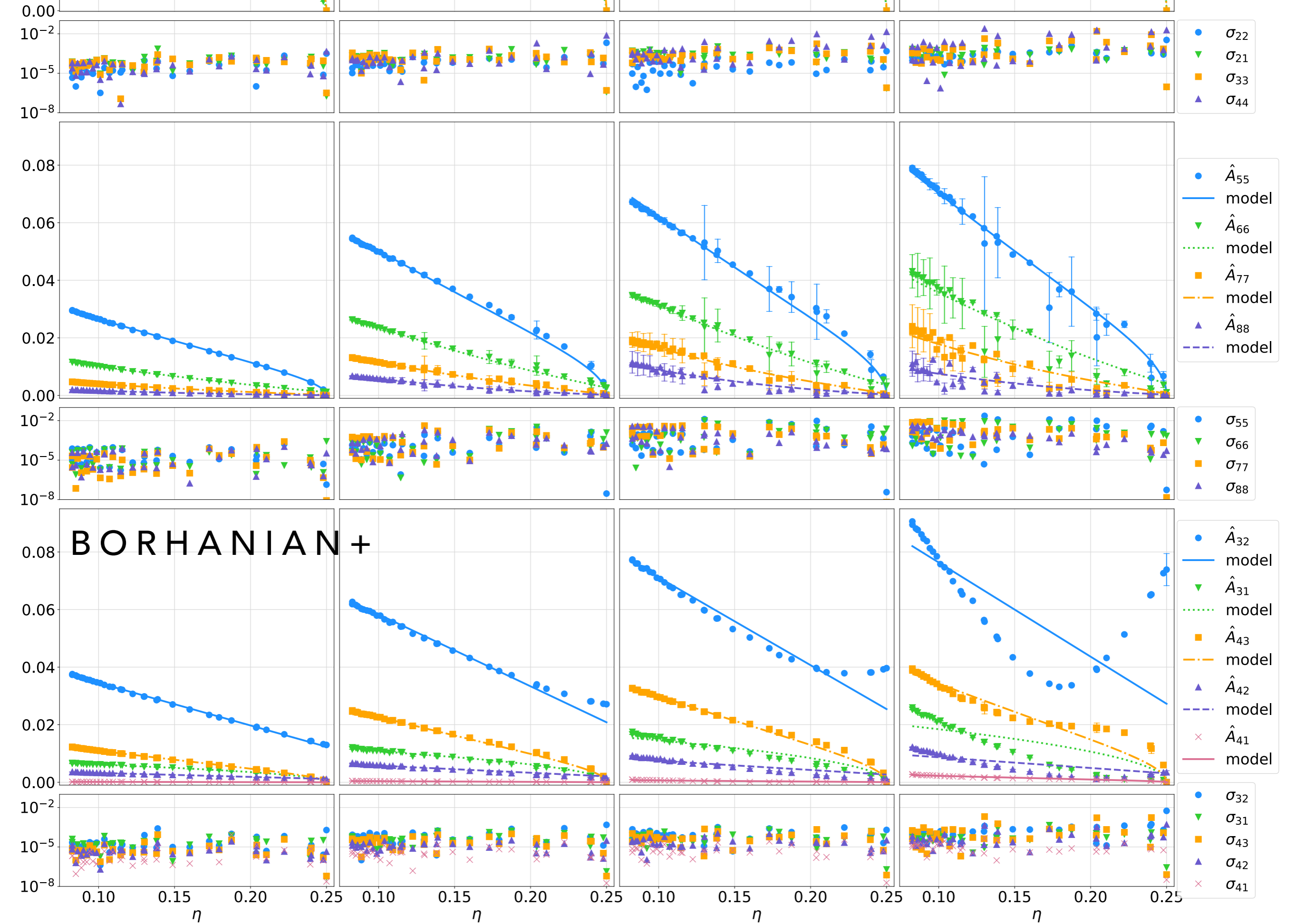
SXS #191, $q=0.40$



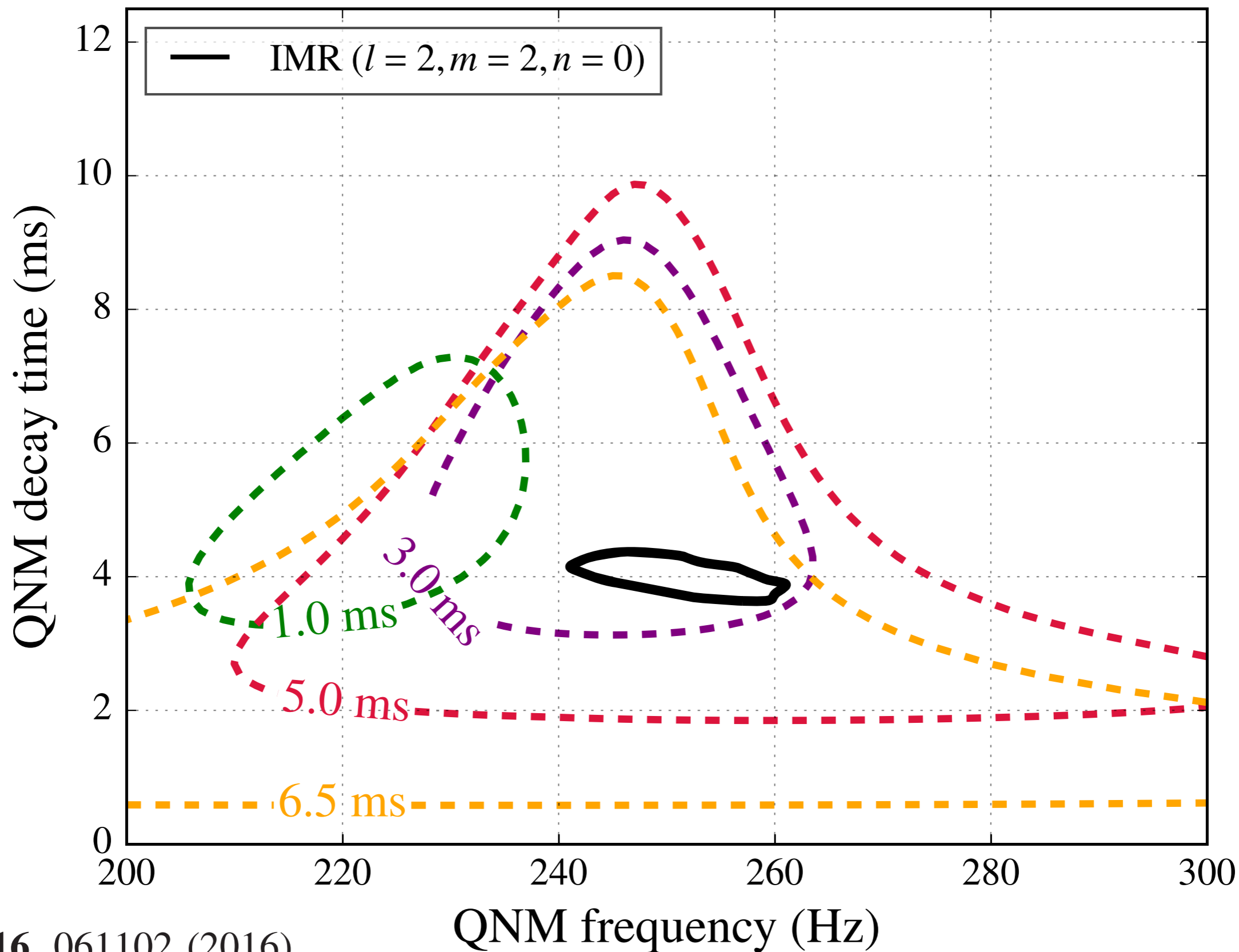
BORHANIAN+

SIGNATURE OF NONLINEARITY IN BINARY BLACK HOLE MERGERS





QUASI-NORMAL MODES IN LIGO'S BLACK HOLES? GW150914



HOW MIGHT ONE TEST THE NO HAIR THEOREM?

- in general relativity the parameters of QNM signal are

$$\vec{\theta}_{\text{GR}} = \{M, \nu, j, \chi_{\text{eff}}, D_{\text{L}}, \theta, \varphi, \psi, \iota, \phi, t_0\},$$

- extra hair:

$$\omega_{lm} = \omega_{lm}^{\text{GR}}(M, J) (1 + \delta\hat{\omega}_{lm}),$$

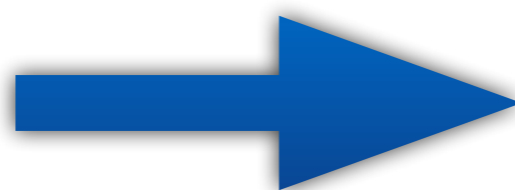
$$\tau_{lm} = \tau_{lm}^{\text{GR}}(M, J) (1 + \delta\hat{\tau}_{lm}),$$

$$H_1 \longleftrightarrow \{\vec{\theta}_{\text{GR}}, \delta\hat{\omega}_{22}\},$$

$$H_2 \longleftrightarrow \{\vec{\theta}_{\text{GR}}, \delta\hat{\omega}_{33}\},$$

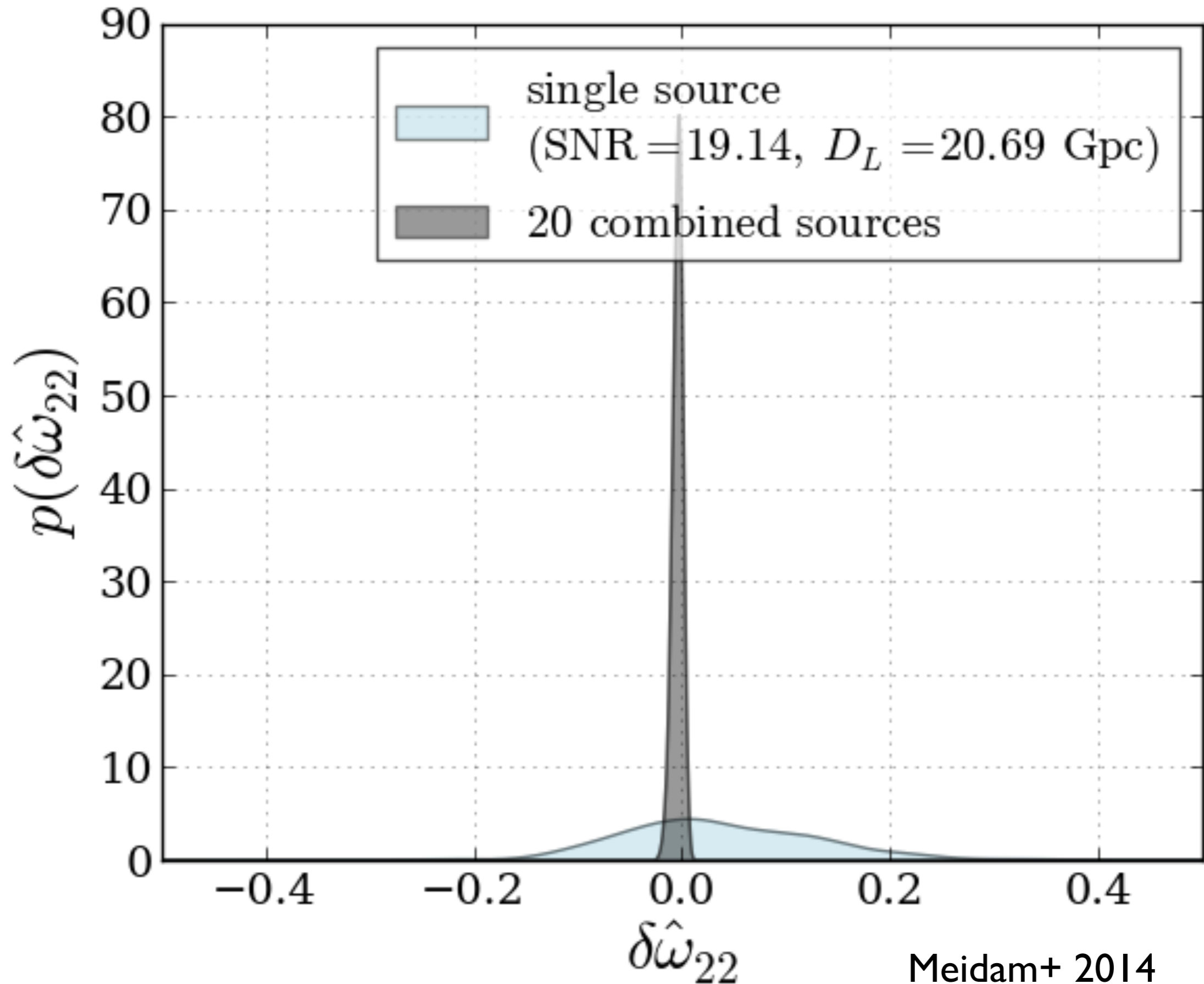
$$H_3 \longleftrightarrow \{\vec{\theta}_{\text{GR}}, \delta\hat{\tau}_{22}\},$$

$$H_{12} \longleftrightarrow \{\vec{\theta}_{\text{GR}}, \delta\hat{\omega}_{22}, \delta\hat{\omega}_{33}\},$$



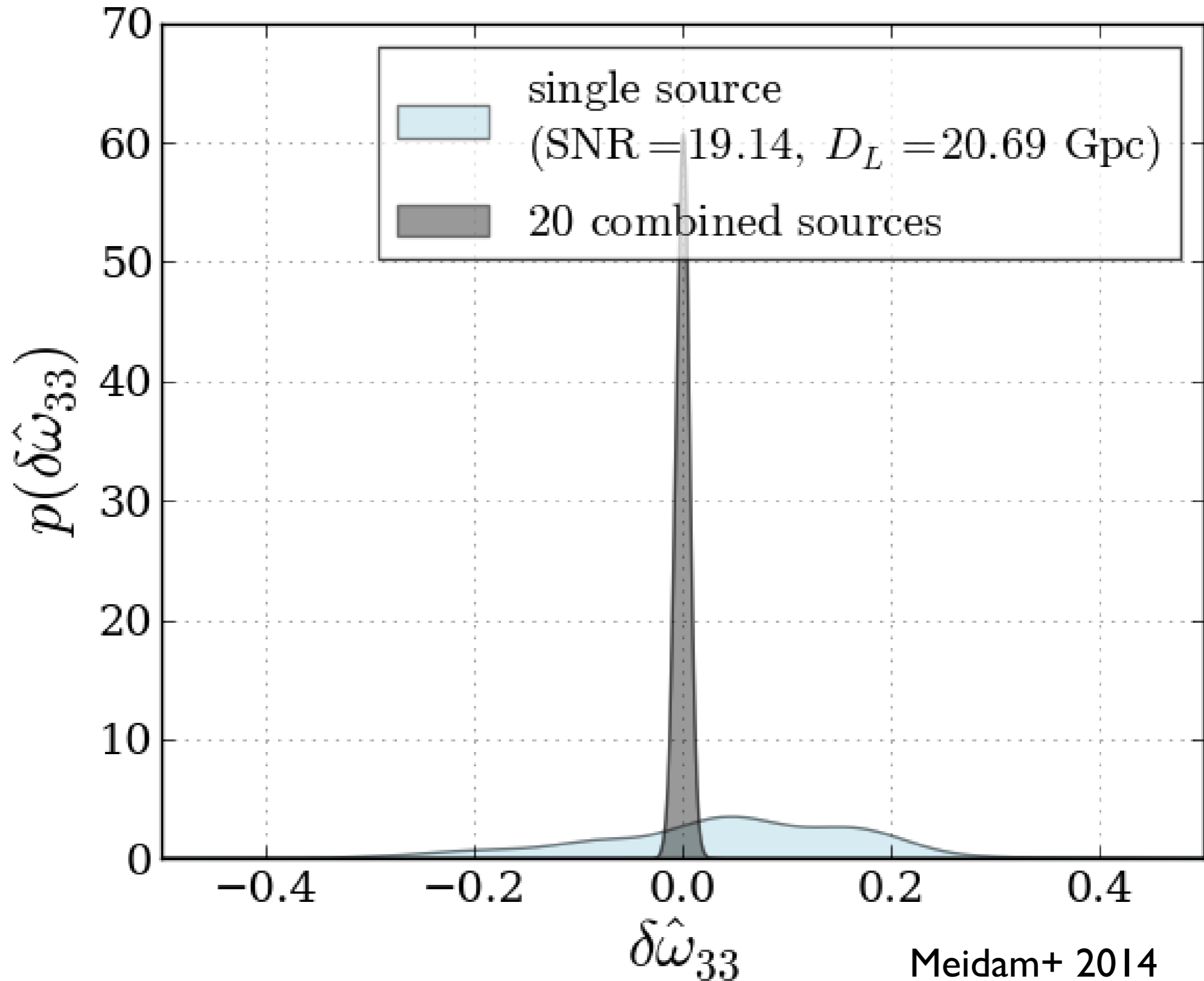
$$B_{\text{GR}}^{123} = \frac{P(d|H_{123}, I)}{P(d|\mathcal{H}_{\text{GR}}, I)}.$$

HOW WELL CAN WE MEASURE NON-GR



Meidam+ 2014

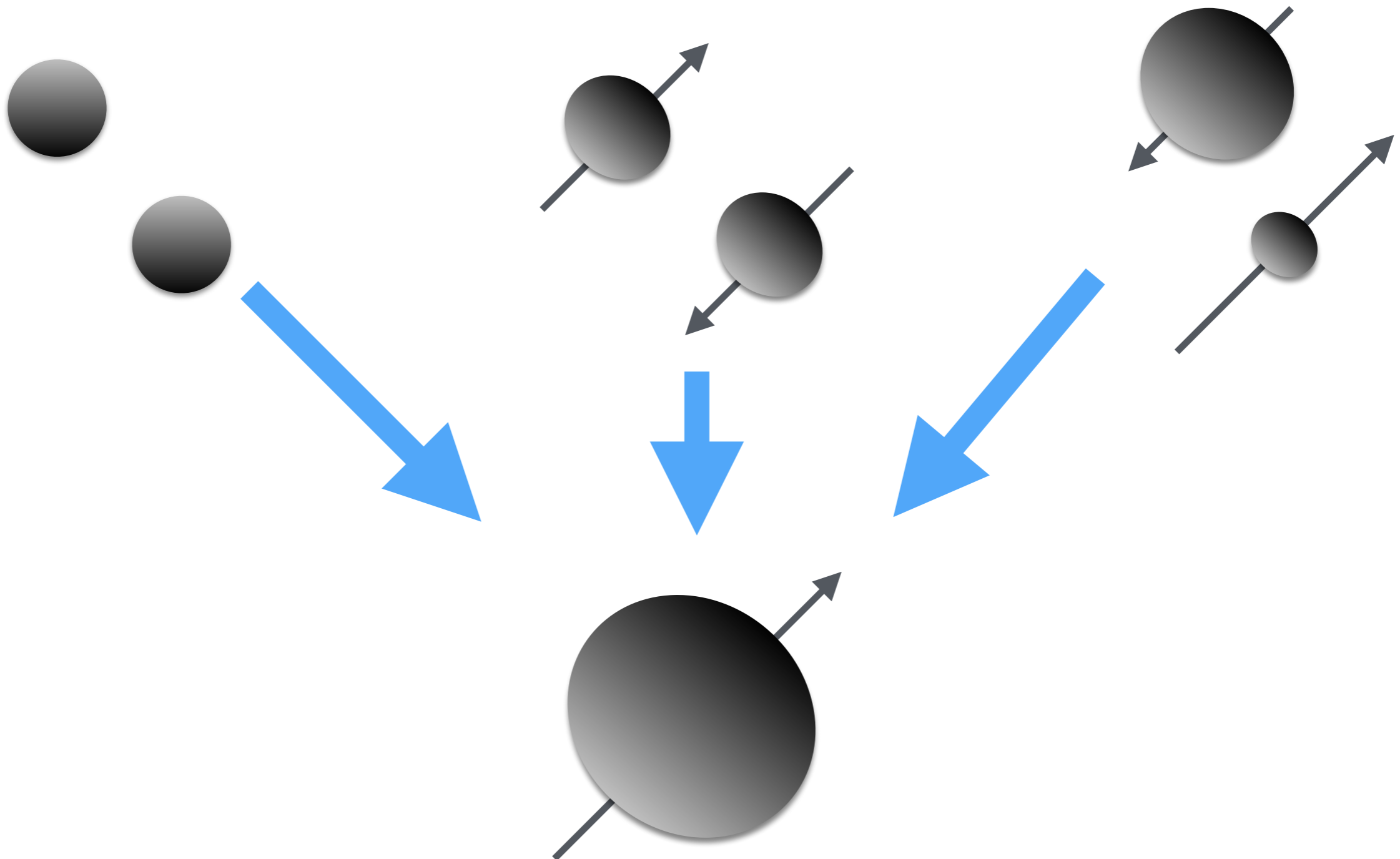
HOW WELL CAN WE MEASURE NON-GR



Meidam+ 2014

NOT ALL BLACK HOLES ARE BORN EQUAL

- infinitely many binaries can form the same final black hole



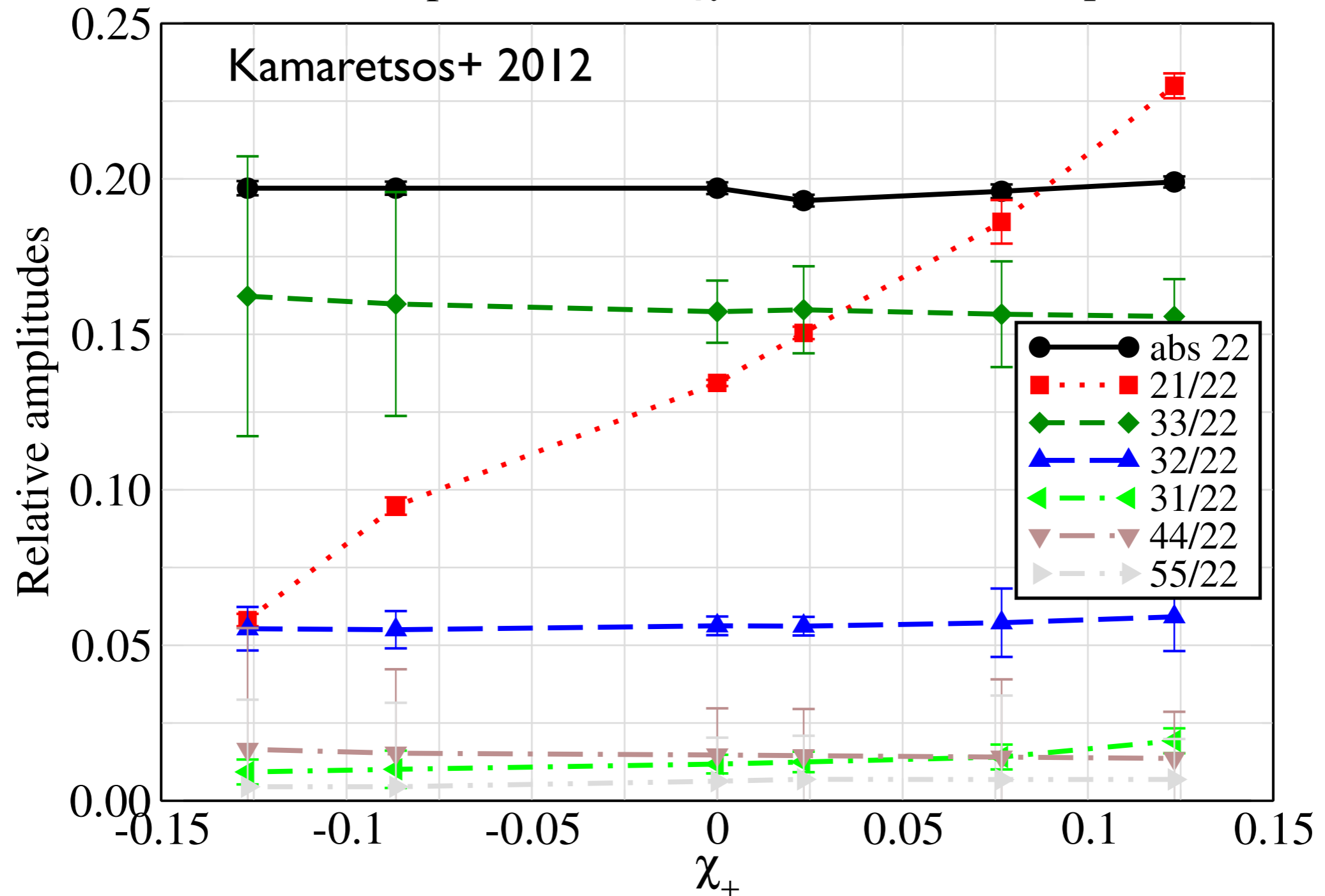
BLACK HOLES NEVER FORGET

22 AND 33 MODE AMPLITUDES ARE CONSTANT

21 MODE VARIES WITH TOTAL SPIN

$$\chi_+ = (m_1 \chi_1 + m_2 \chi_2) / M_{in}$$

Initial spins such that $\chi \sim 0.62$, mass ratio $q=2$



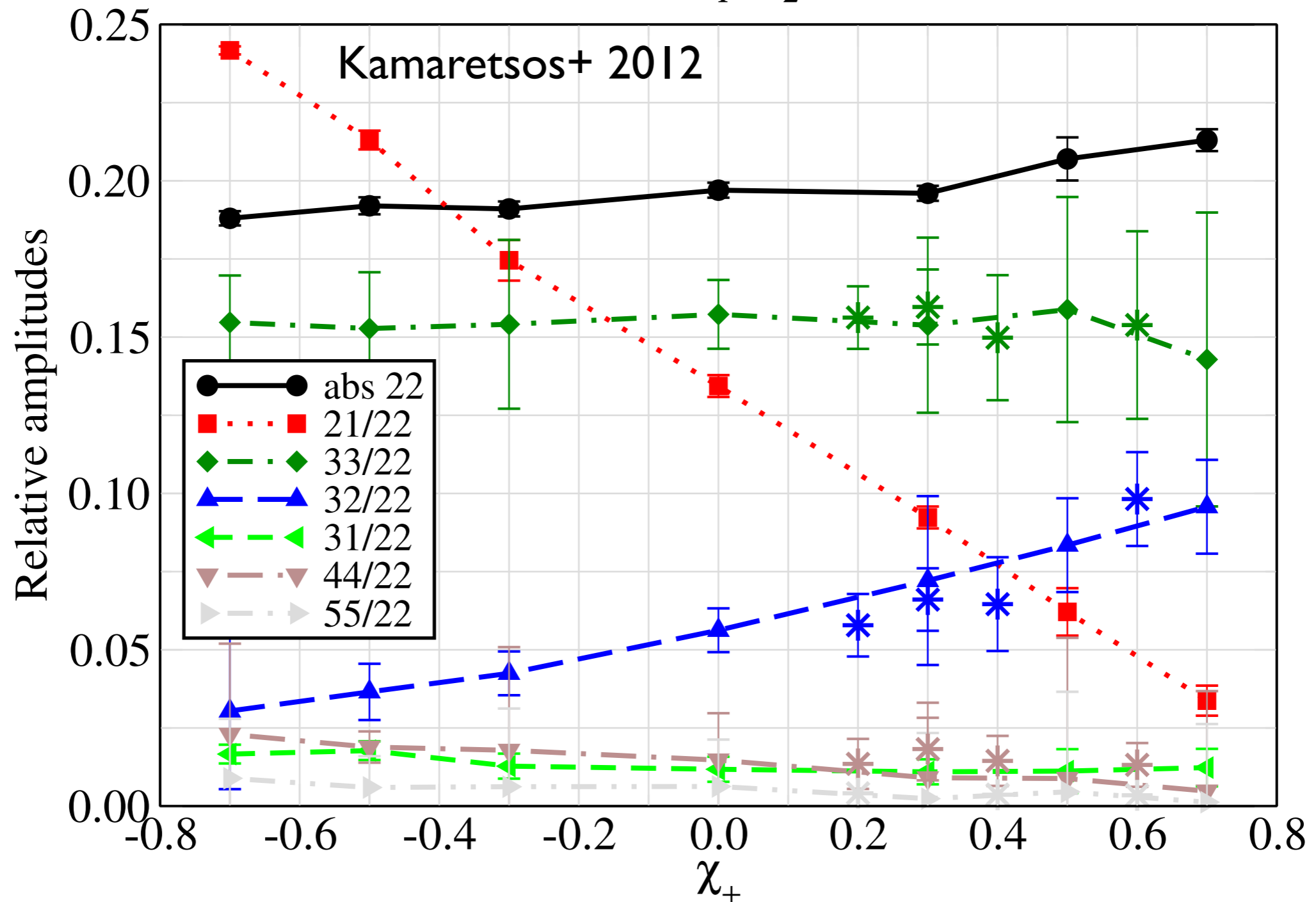
BLACK HOLES NEVER FORGET

22 AND 33 MODE AMPLITUDES ARE CONSTANT

21 MODE VARIES WITH TOTAL SPIN

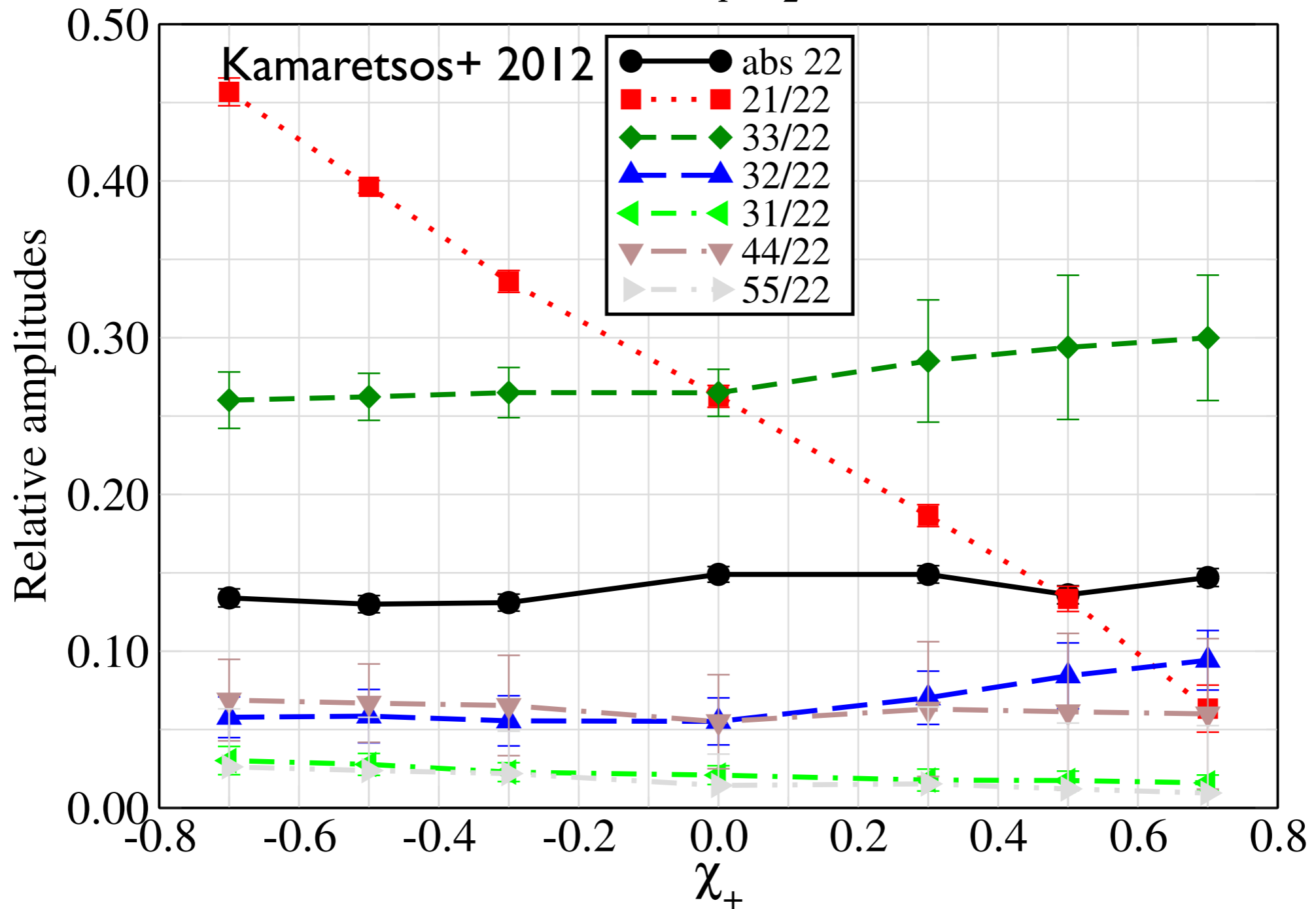
$$\chi_+ = (m_1 \chi_1 + m_2 \chi_2) / M_{in}$$

Equal initial spins $\chi_1 = \chi_2$, mass ratio $q=2$

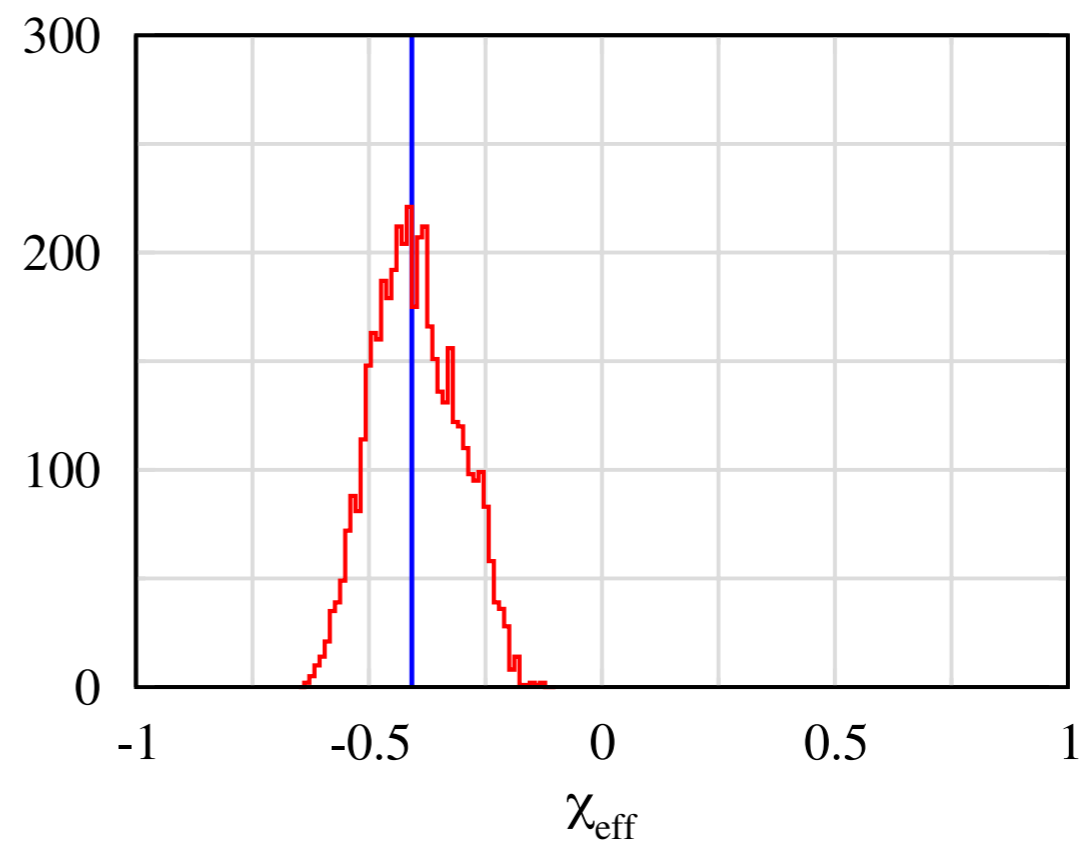
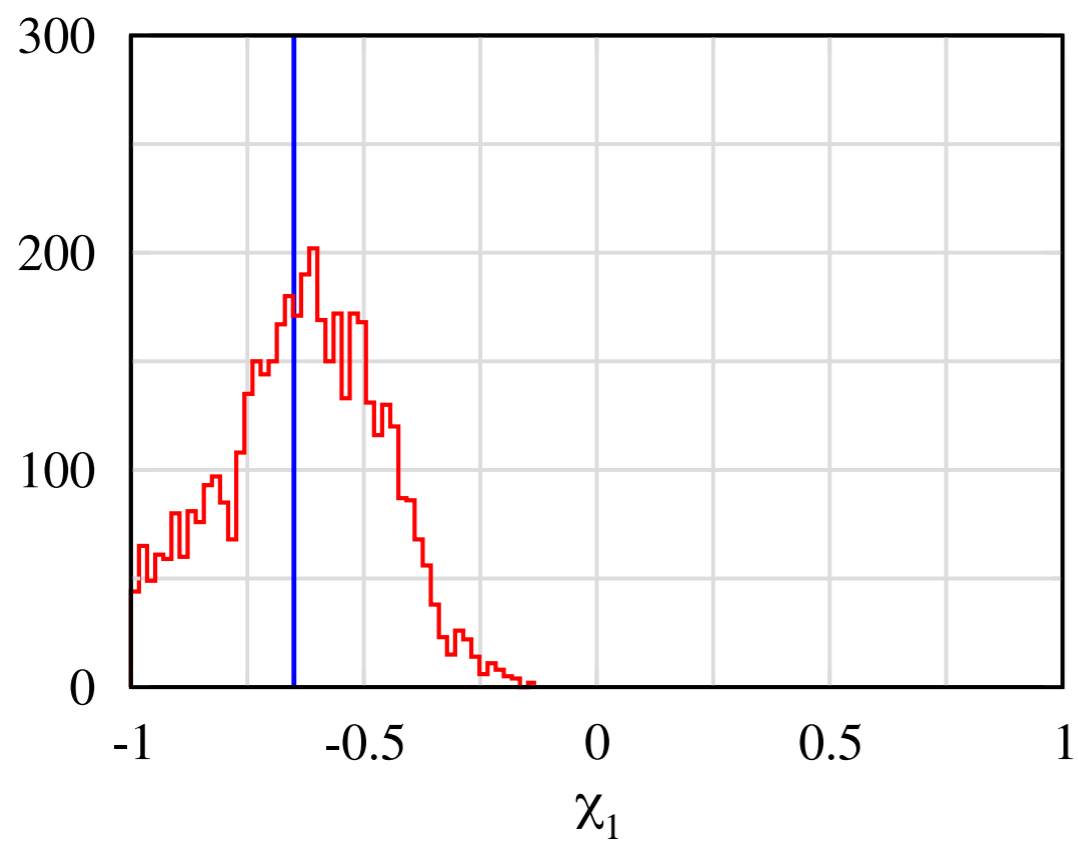
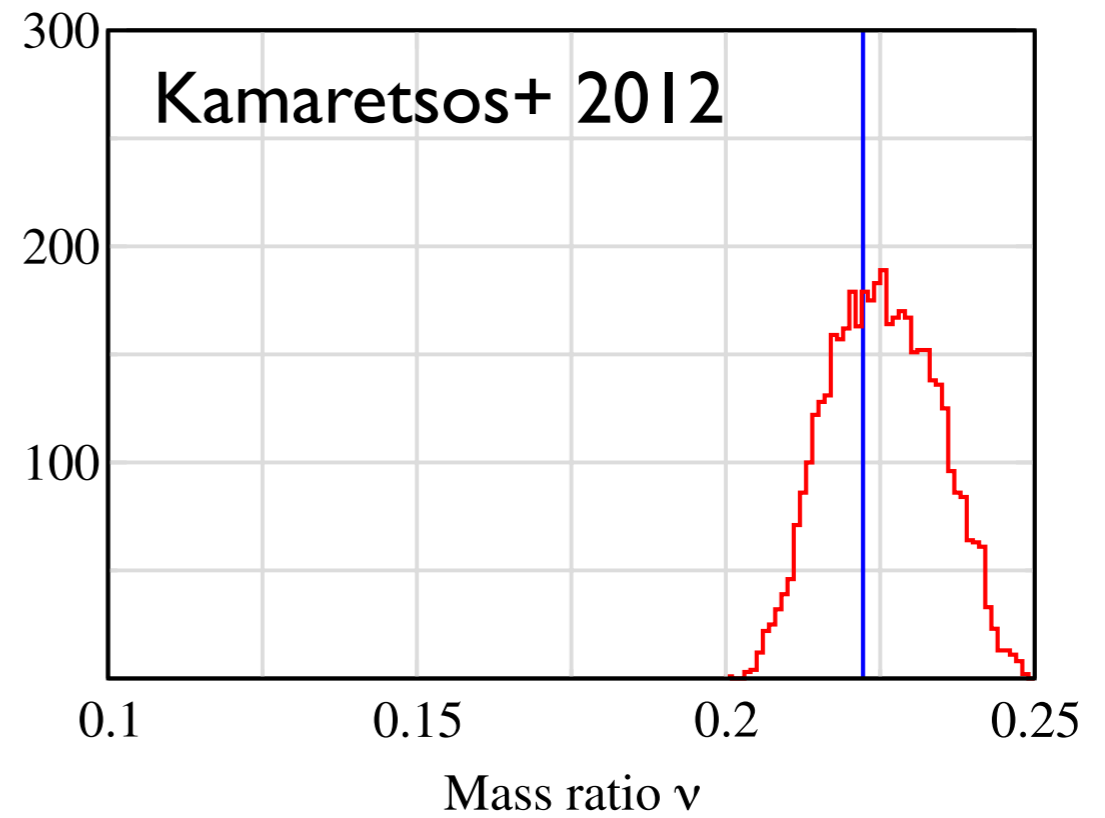
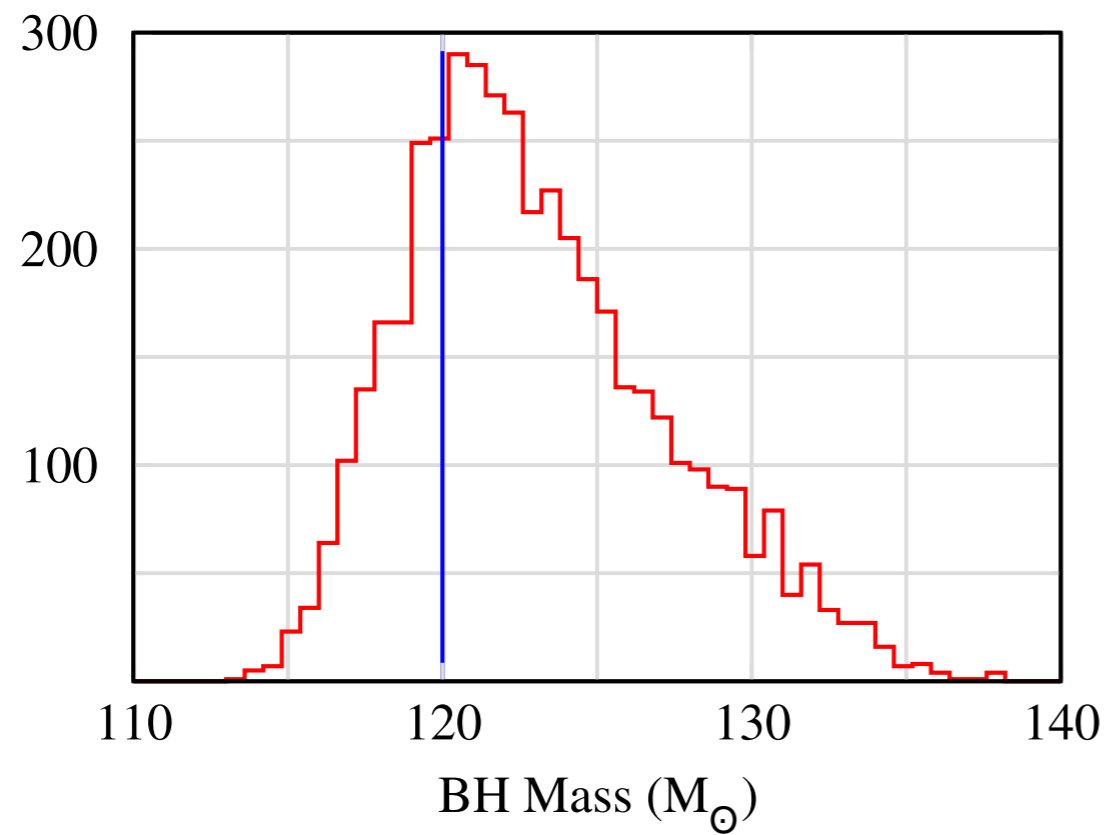


IS THIS GENERIC?

Equal initial spins $\chi_1 = \chi_2$, mass ratio $q=4$



FROM RINGDOWN BACK TO INSPIRAL

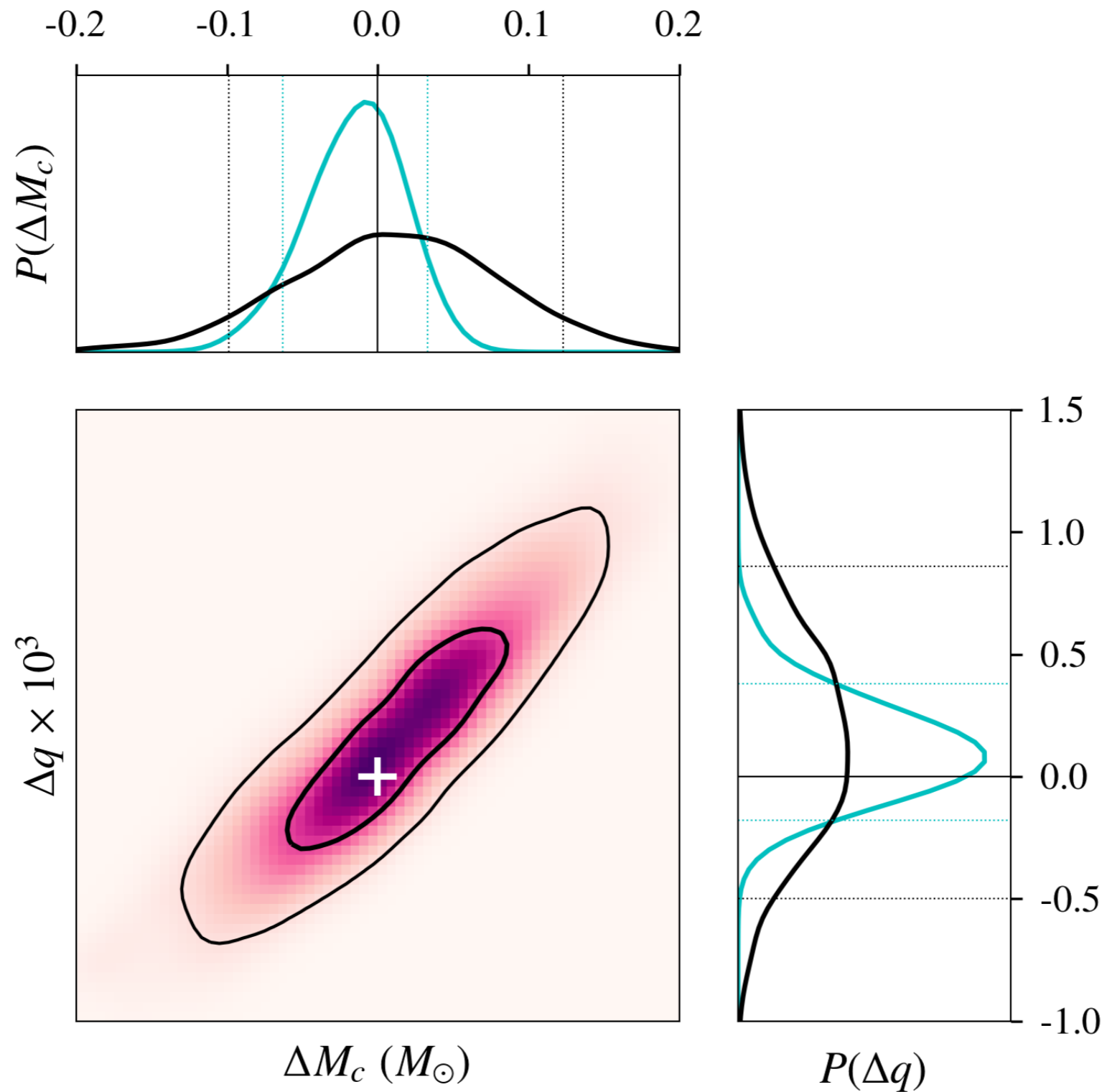


HIGHER HARMONICS IN BLACK HOLE BINARIES

$$h(t; \mathbf{n}, \lambda) = \frac{1}{d_L} \sum_{\ell=2}^{\infty} \sum_{m=-\ell}^{\ell} Y_{\ell m}^{-2}(\mathbf{n}) h_{\ell m}(t; \lambda),$$

$$h(t; \mathbf{n}, \lambda, \Delta\lambda) = \sum_{m=\pm 2} Y_{2m}^{-2}(\mathbf{n}) h_{2m}(t, \lambda) \\ + \sum_{\text{H.O.M}} Y_{\ell m}^{-2}(\mathbf{n}) h_{\ell m}(t, \lambda + \Delta\lambda)$$

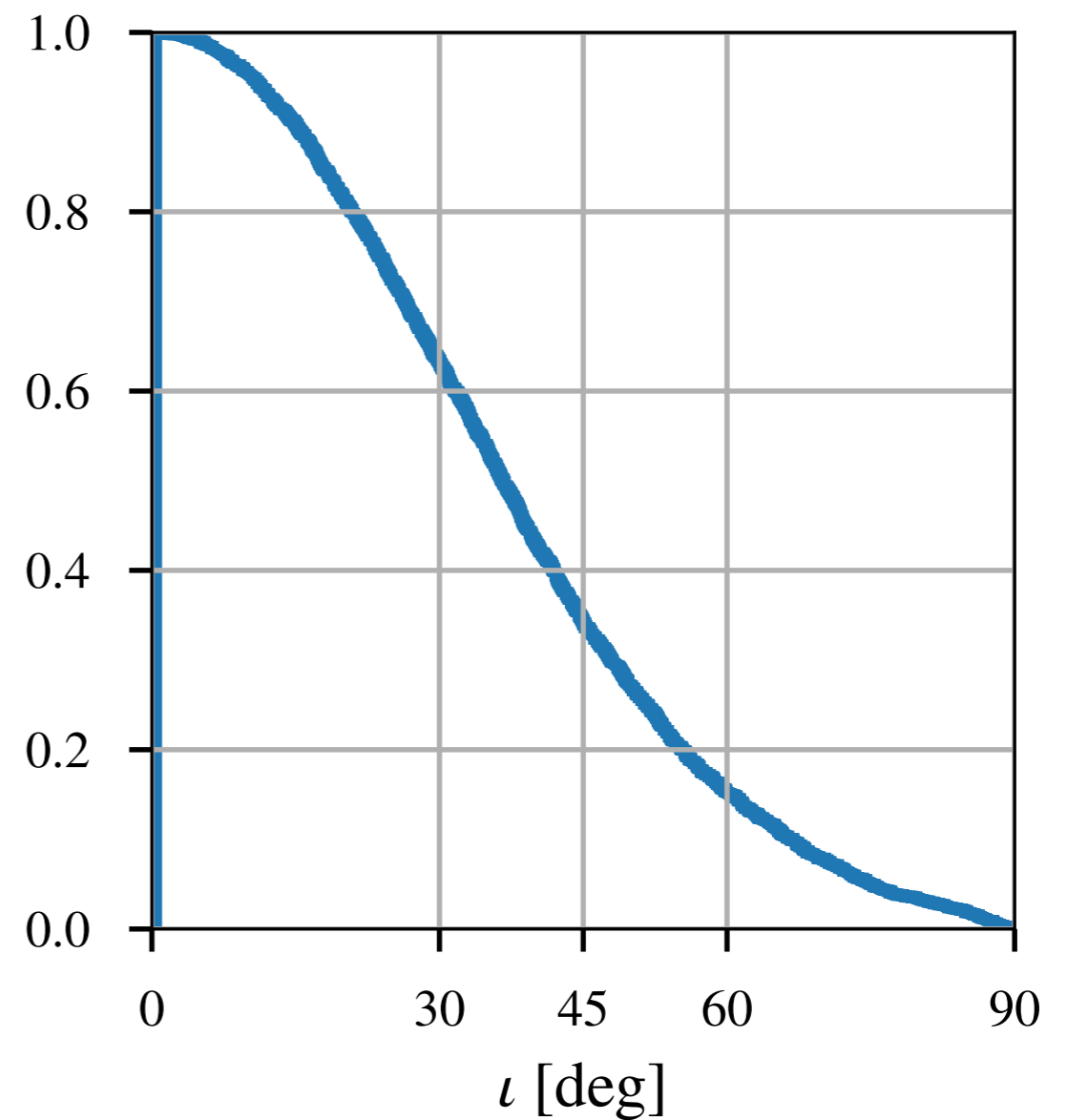
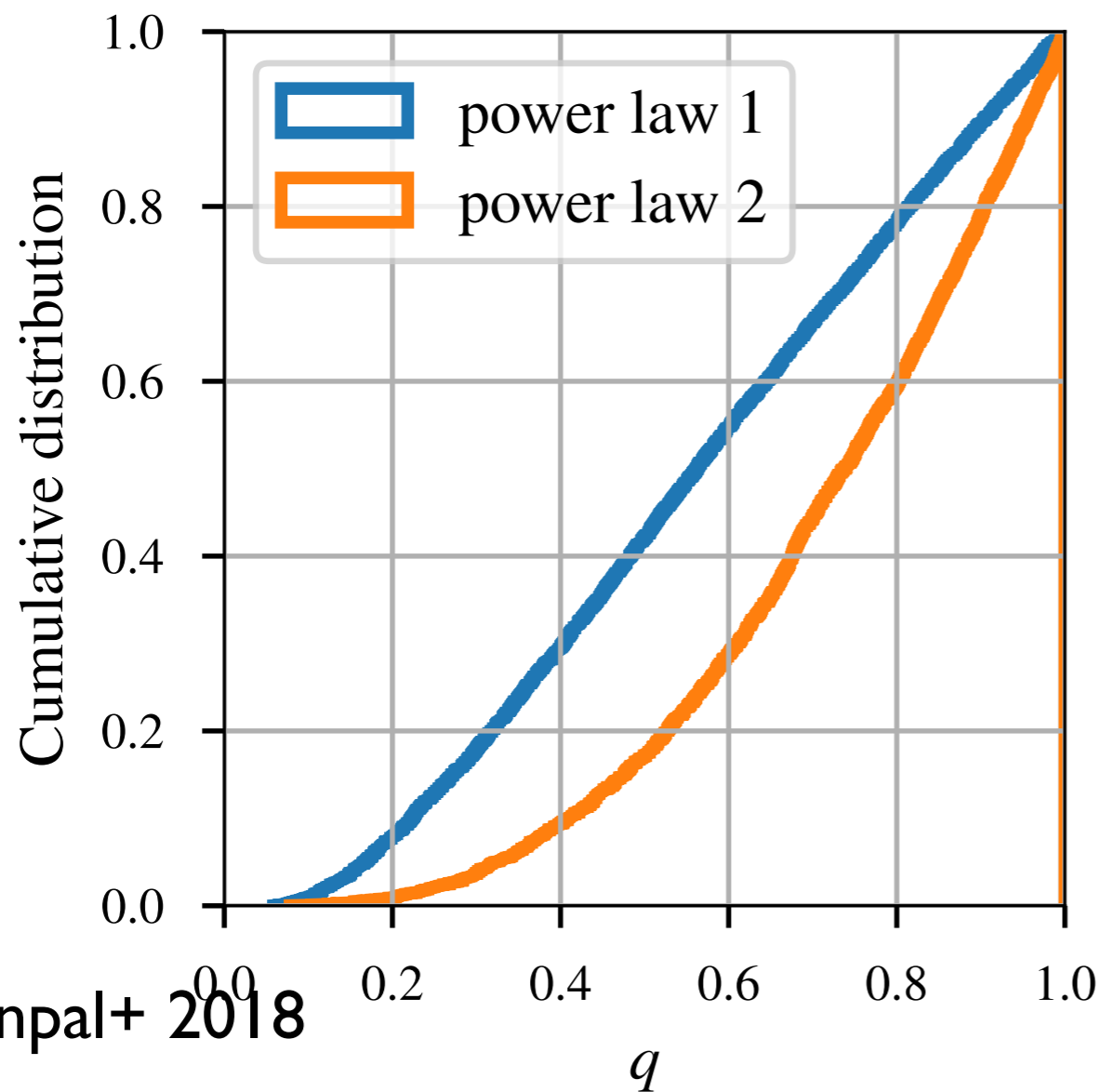
A "NO-HAIR" TEST FOR BINARY BLACK HOLES



Dhanpal+ 2018

CHANCE OF DETECTING HIGHER MODES?

- test is effective only when the mass ratio is small (say $< 1/2$) and inclination is large (say > 60 degrees)
- only a few percent of all detected sources are suitable



WORK DONE IN COLLABORATION WITH

- LIGO Scientific Collaboration
- Virgo Collaboration
- **Agathos**, Ajith, Arun, **Bohranian**, **Dhanpal**, Hannam, Husa, **Kamaretsos**, **Ghosh**, **Gossan**, **Meidam**, **Mehta**, van den Broeck, Veitch

